



Chemical Propulsion Systems

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- Introduction
- Fundamental equations
- Thermodynamics of gases
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- Isentropic flow
- Nozzle fluid flow



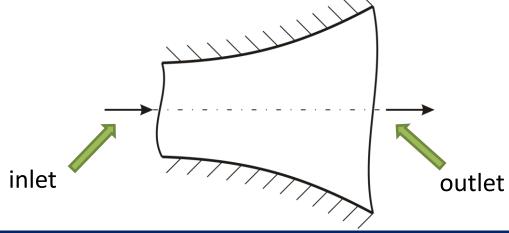
Introduction

Fluid flow:

- naturally three dimensional, but in some special cases can be considered as one dimensional or quasi-one dimension
- fluid can be considered according to:
 - steady-state VS transient
 - turbulent VS laminar

inviscid VS viscous fluid

quasi-one dimension fluid flow

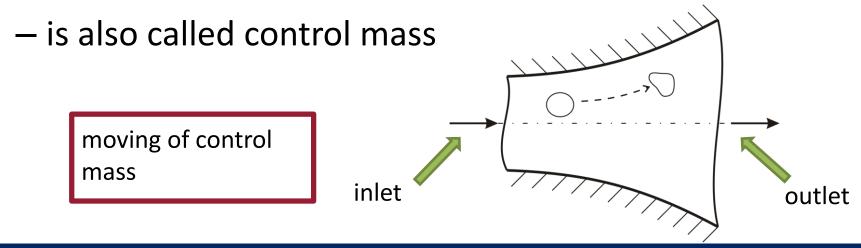






Introduction

- System or Control Mass (CM):
 - is a collection of matter of fixed identity
 - it may be considered enclosed by an invisible, massless, flexible surface through which no matter can pass
 - the boundary of the system may change position, size, and shape

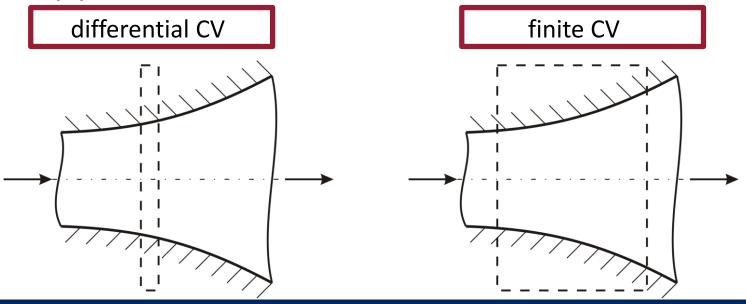






Introduction

- Control Volume (CV):
 - is arbitrary volume fixed to the coordinate system (stationary or moving)
 - bounded by control surface (CS) through which fluid may pass, CV can has differential or finite size





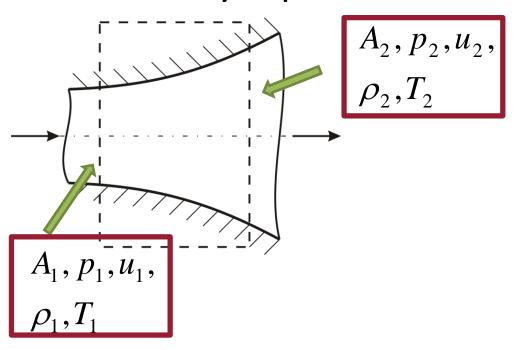


Fundamental equations

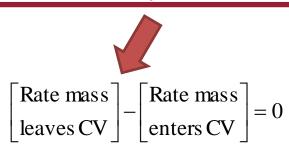
- There are 4 fundamental equations, which must be considered:
 - Continuity equation
 - Momentum equation
 - Energy equation
 - Entropy equation



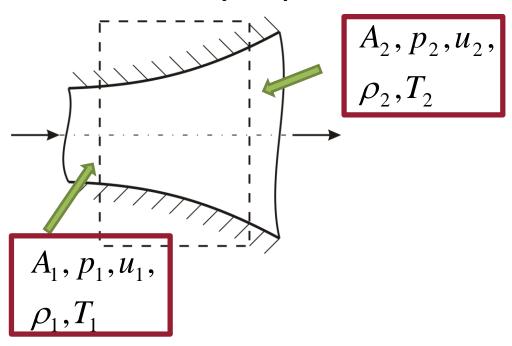
Continuity equation



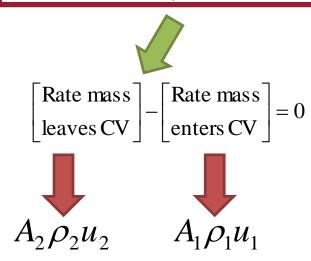
Fundamental equations



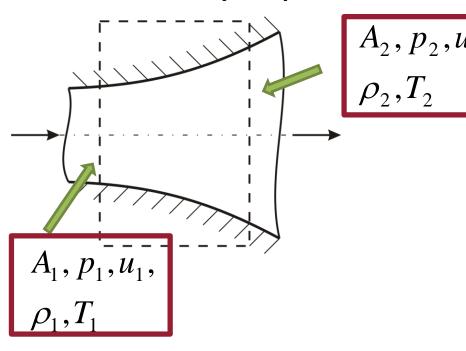
Continuity equation



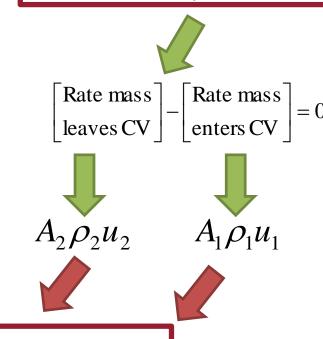
Fundamental equations



Continuity equation



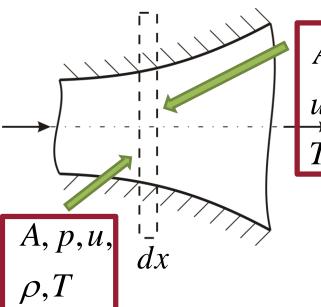
Fundamental equations



$$A_2 \rho_2 u_2 - A_1 \rho_1 u_1 = 0$$

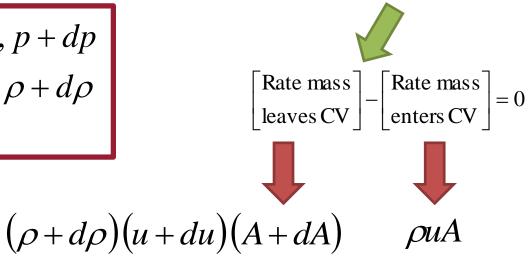
Fundamental equations

Continuity equation



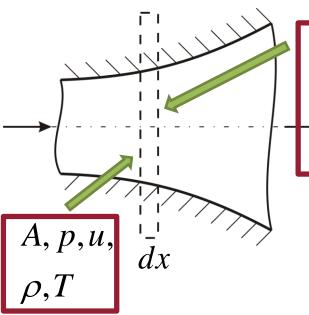
$$A + dA, p + dp$$

 $u + du, \rho + d\rho$
 $T + dT$



Fundamental equations

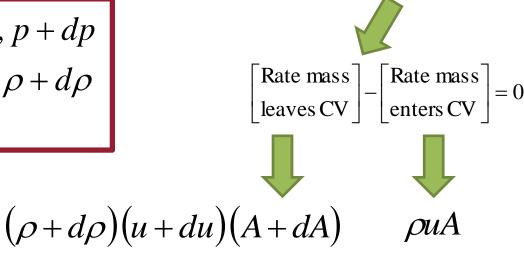
Continuity equation



$$A + dA, p + dp$$

$$u + du, \rho + d\rho$$

$$T + dT$$

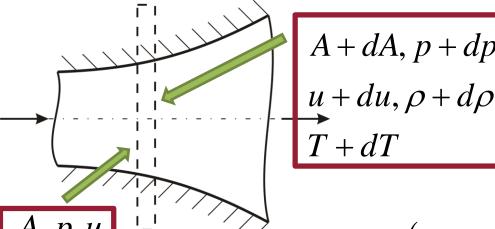




$$\rho udA + \rho Adu + uAd\rho = 0$$

Fundamental equations

Continuity equation



A + dA, p + dp $u + du, \rho + d\rho$







$$(\rho + d\rho)(u + du)(A + dA)$$





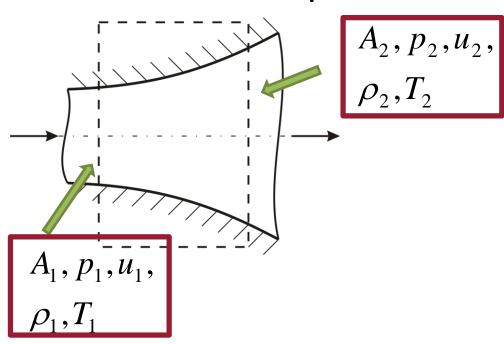
$$\frac{d\rho}{\rho} + \frac{du}{u} + \frac{dA}{A} = 0$$

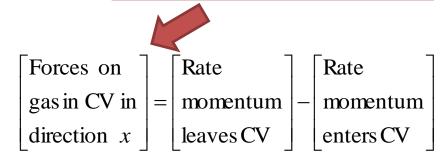


$$\rho udA + \rho Adu + uAd\rho = 0$$

Fundamental equations

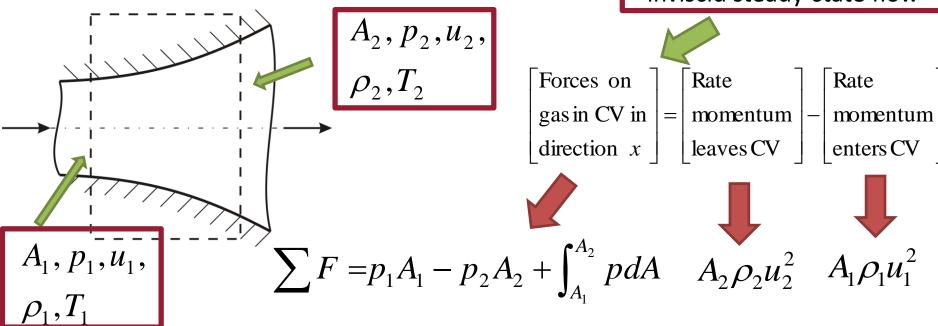
Momentum equation





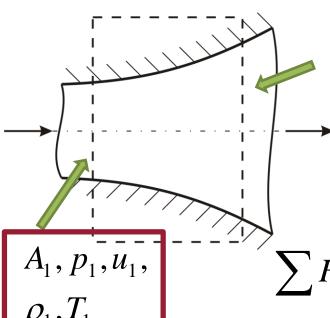
Fundamental equations

Momentum equation



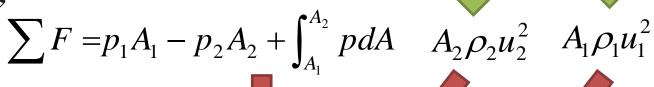
Fundamental equations

Momentum equation



$$A_2, p_2, u_2, \\ \rho_2, T_2$$

$$\begin{bmatrix} \text{Forces on} \\ \text{gas in CV in} \\ \text{direction } x \end{bmatrix} = \begin{bmatrix} \text{Rate} \\ \text{momentum} \\ \text{leaves CV} \end{bmatrix} - \begin{bmatrix} \text{Rate} \\ \text{momentum} \\ \text{enters CV} \end{bmatrix}$$









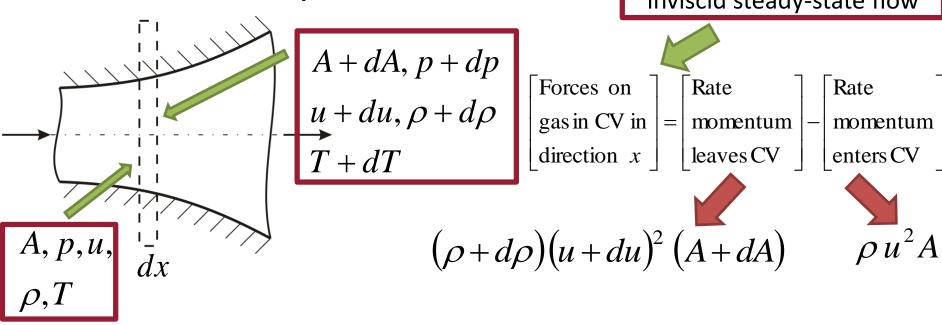


$$p_1 A_1 - p_2 A_2 + \int_{A_1}^{A_2} p dA = A_2 \rho_2 u_2^2 - A_1 \rho_1 u_1^2$$



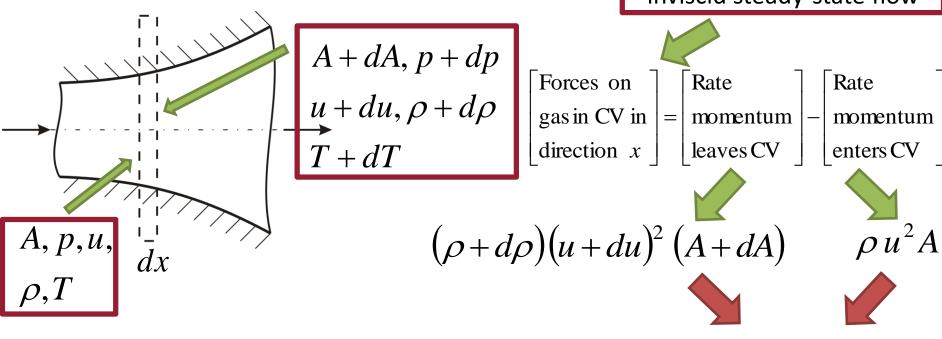
Fundamental equations

Momentum equation



Fundamental equations

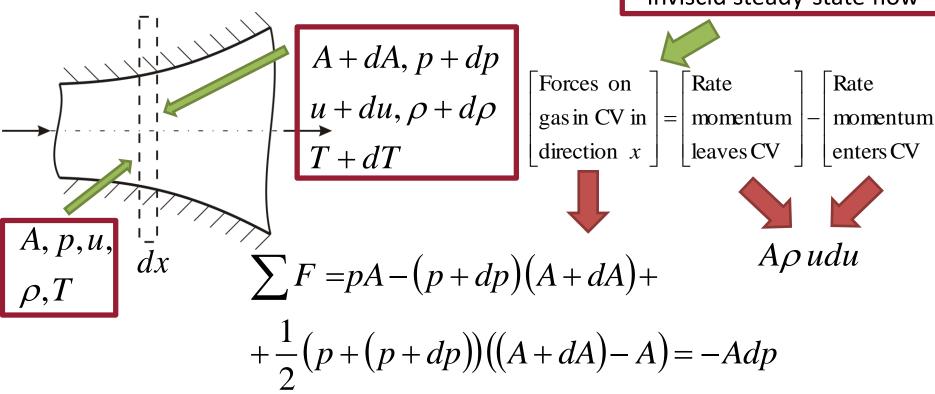
Momentum equation





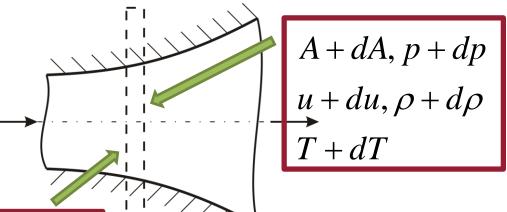
Fundamental equations

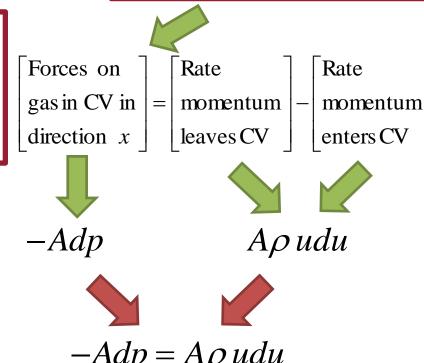
Momentum equation



Fundamental equations

Momentum equation





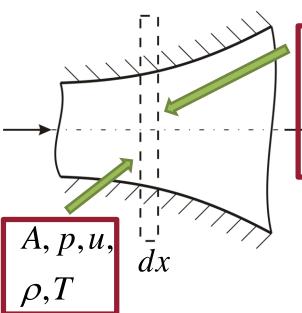
$$-Adp = A\rho udu$$



Fundamental equations

Momentum equation

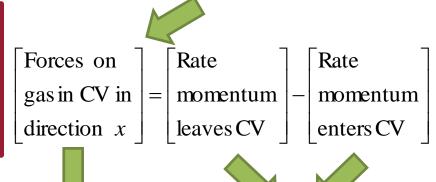
inviscid steady-state flow



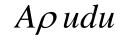
$$A + dA, p + dp$$

$$u + du, \rho + d\rho$$

$$T + dT$$



$$-Adp$$







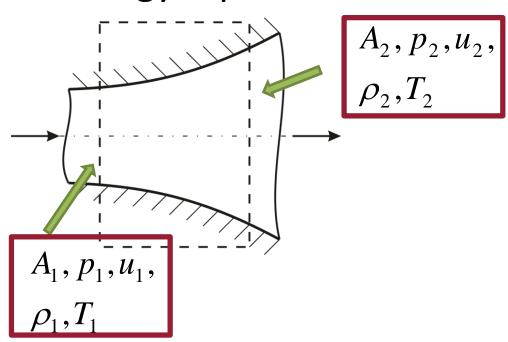
$$-\frac{dp}{\rho} = udu$$

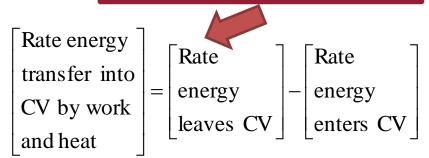
$$-Adp = A\rho udu$$

Euler's equation

Fundamental equations

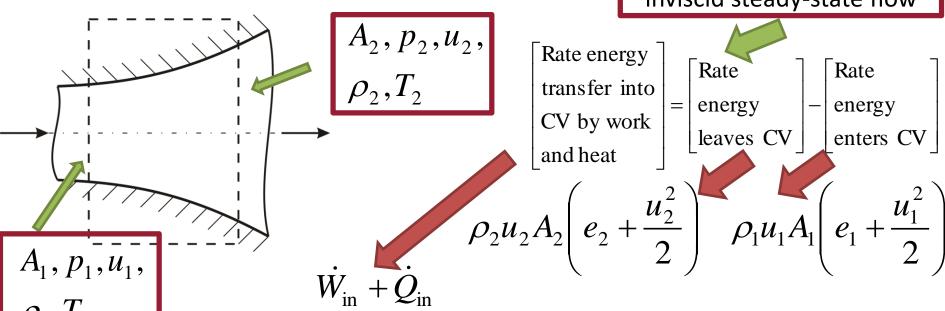
Energy equation





Fundamental equations

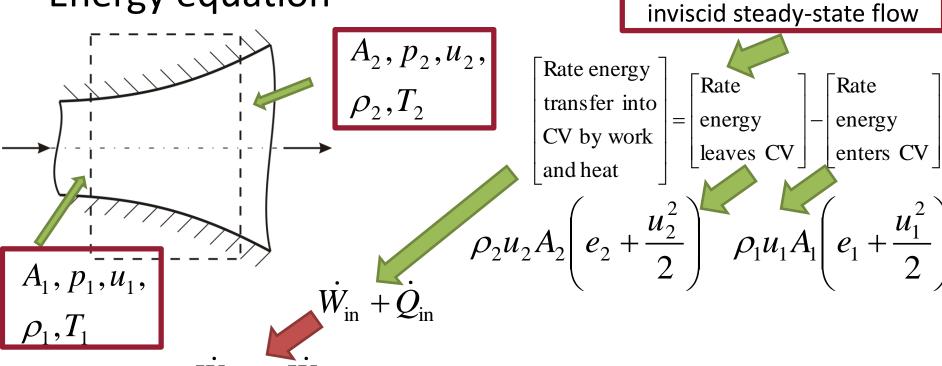
Energy equation





Fundamental equations

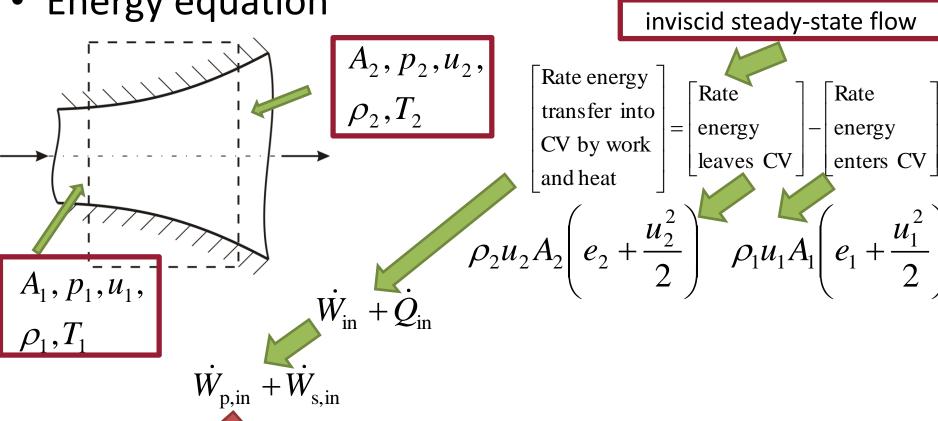
Energy equation





Fundamental equations

Energy equation



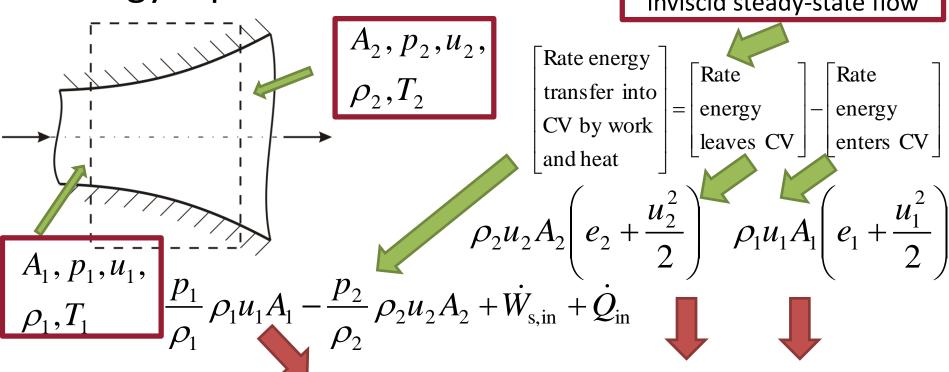
 $\frac{p_1}{\rho_1} \rho_1 u_1 A_1 - \frac{p_2}{\rho_2} \rho_2 u_2 A_2$





Fundamental equations

Energy equation

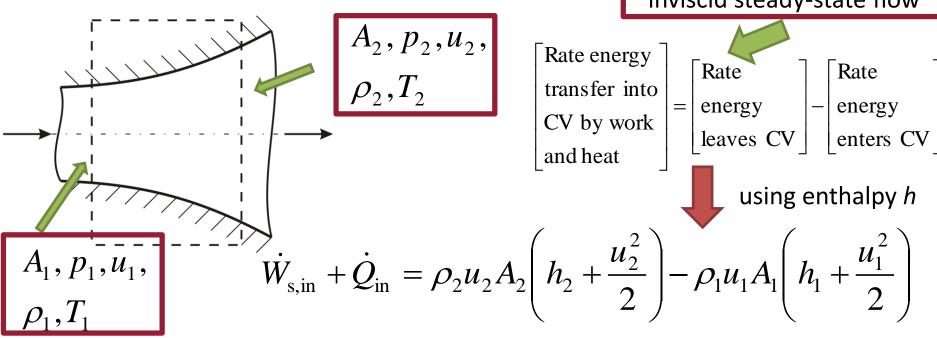


$$\dot{W}_{\text{s,in}} + \dot{Q}_{\text{in}} = \rho_2 u_2 A_2 \left(e_2 + \frac{p_2}{\rho_2} + \frac{u_2^2}{2} \right) - \rho_1 u_1 A_1 \left(e_1 + \frac{p_1}{\rho_1} + \frac{u_1^2}{2} \right)$$



Fundamental equations

Energy equation

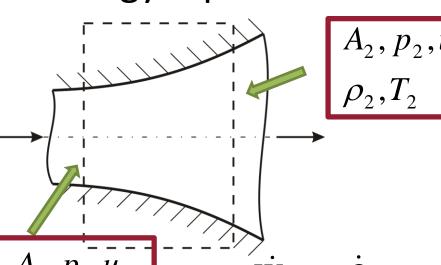




Fundamental equations

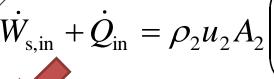
Energy equation

inviscid steady-state flow



$$\begin{bmatrix} \text{Rate} \\ \text{energy} \\ \text{leaves CV} \end{bmatrix} - \begin{bmatrix} \text{Rate} \\ \text{energy} \\ \text{enters CV} \end{bmatrix}$$

using enthalpy h



$$\dot{W}_{\text{s,in}} + \dot{Q}_{\text{in}} = \rho_2 u_2 A_2 \left(h_2 + \frac{u_2^2}{2} \right) - \rho_1 u_1 A_1 \left(h_1 + \frac{u_1^2}{2} \right)$$

adiabatic process no shaft work



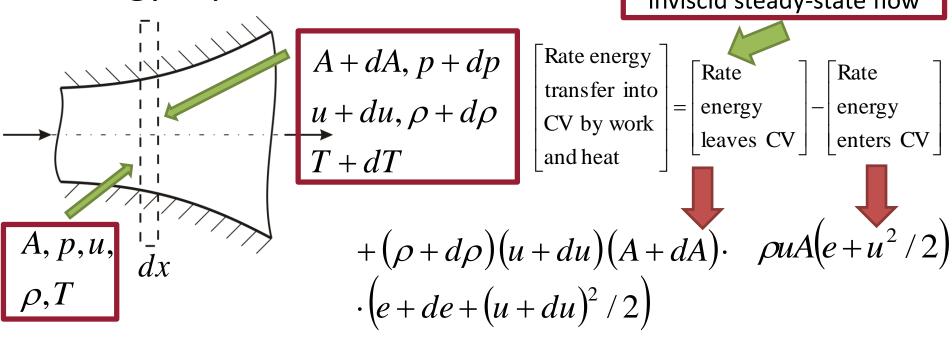
$$\left(h_1 + \frac{u_1^2}{2}\right) = \left(h_2 + \frac{u_2^2}{2}\right)$$





Fundamental equations

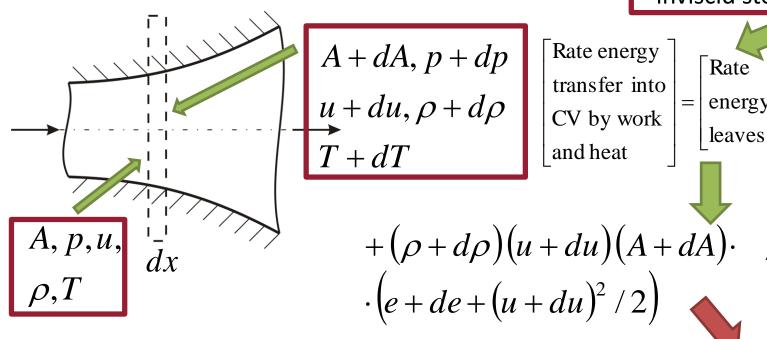
Energy equation

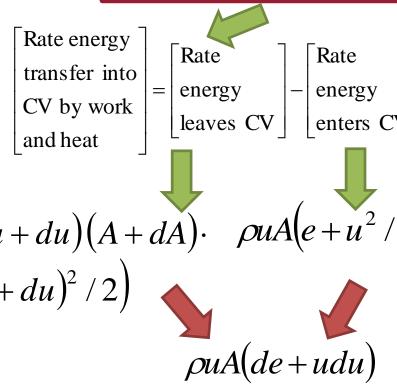




Fundamental equations

Energy equation

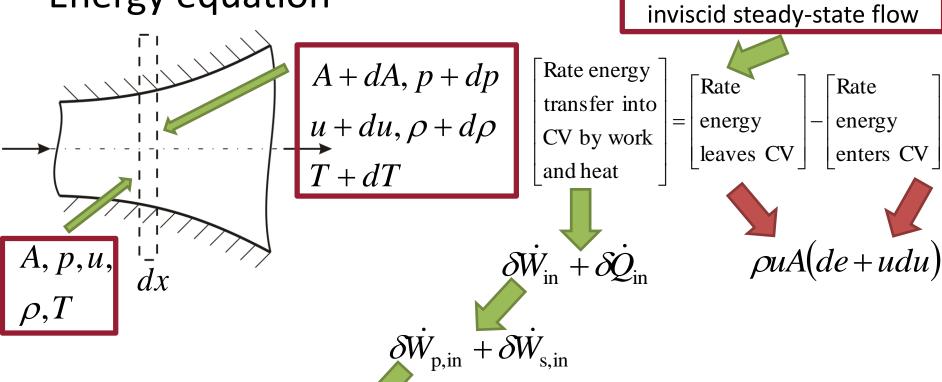






Fundamental equations

Energy equation



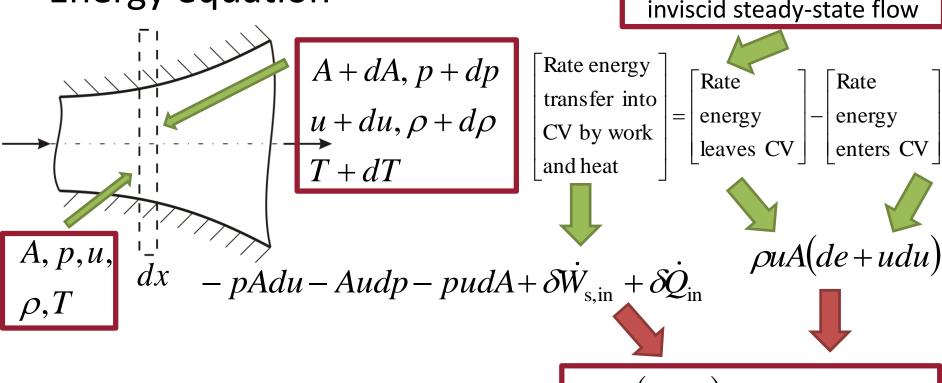
$$puA - (p+dp)(u+du)(A+dA)$$
 \longrightarrow $-pAdu-Audp-pudA$





Fundamental equations

Energy equation



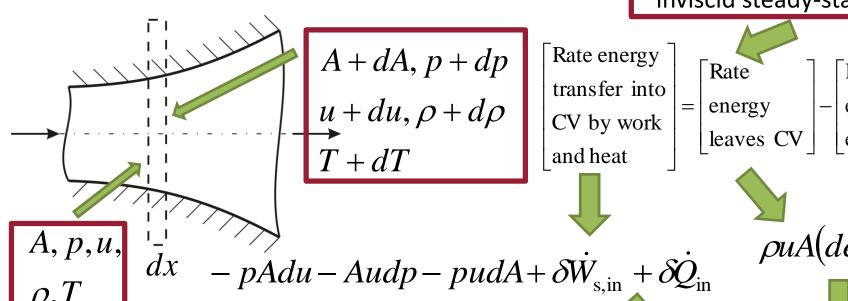
$$-d(p/\rho) + \delta w_{s,in} + \delta q_{in} =$$

$$= de + udu$$

Fundamental equations

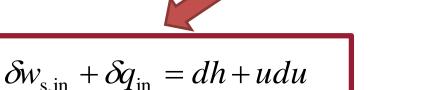
Energy equation

inviscid steady-state flow



 $\rho uA(de+udu)$

using enthalpy h





$$-d(p/\rho) + \delta w_{s,in} + \delta q_{in} =$$

$$= de + udu$$

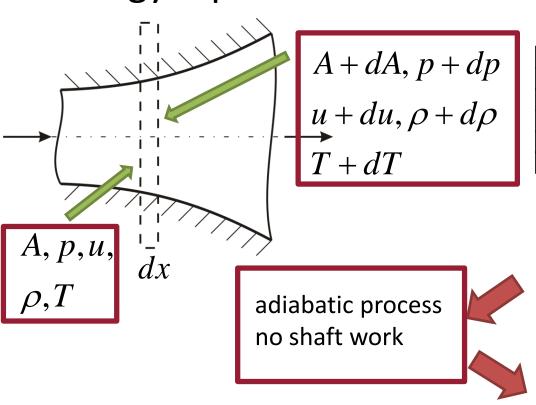


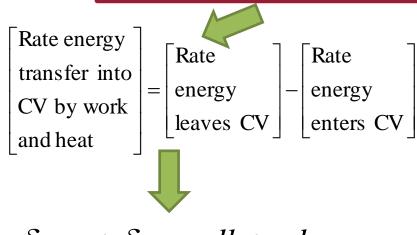


Fundamental equations

Energy equation



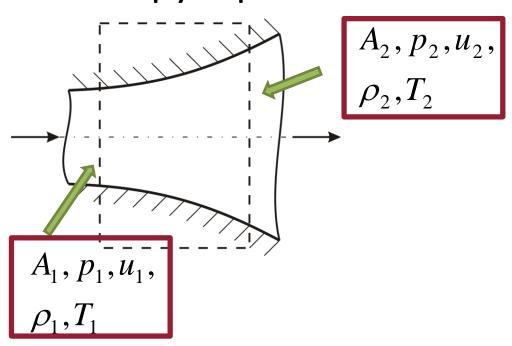




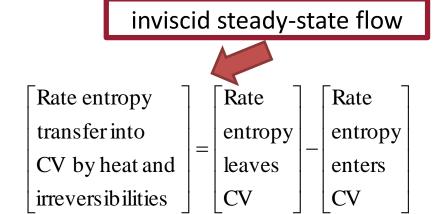
$$\delta w_{\rm s,in} + \delta q_{\rm in} = dh + udu$$

$$dh + udu = 0$$

Entropy equation

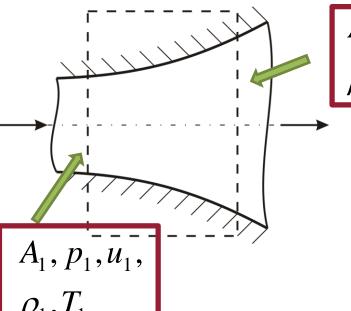


Fundamental equations

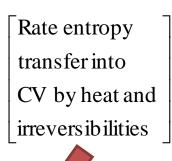


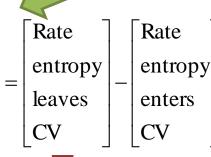
Fundamental equations

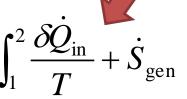
Entropy equation



$$\begin{bmatrix} A_2, p_2, u_2, \\ \rho_2, T_2 \end{bmatrix}$$





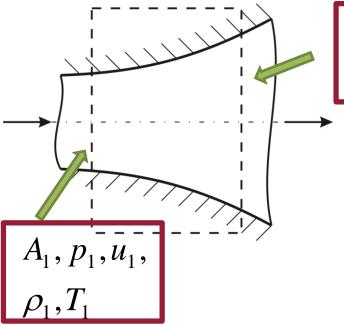






Fundamental equations

Entropy equation



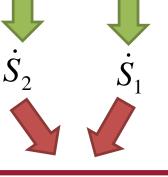
$$A_2, p_2, u_2,$$

$$\rho_2, T_2$$

inviscid steady-state flow

$$\begin{bmatrix} \text{Rate entropy} \\ \text{trans fer into} \\ \text{CV by heat and} \\ \text{irreversibilities} \end{bmatrix} = \begin{bmatrix} \text{Rate} \\ \text{entropy} \\ \text{leaves} \\ \text{CV} \end{bmatrix} - \begin{bmatrix} \text{Rate} \\ \text{entropy} \\ \text{enters} \\ \text{CV} \end{bmatrix}$$

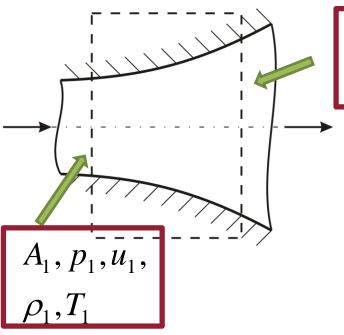
$$\int_{1}^{2} \frac{\delta \dot{Q}_{\rm in}}{T} + \dot{S}_{\rm gen}$$



$$\int_{1}^{2} \frac{\delta \dot{Q}_{\rm in}}{T} + \dot{S}_{\rm gen} = \dot{S}_{2} - \dot{S}_{1}$$

Fundamental equations

Entropy equation

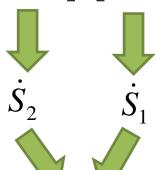


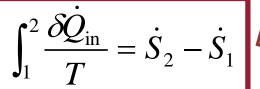
$$A_2, p_2, u_2,$$

$$\rho_2, T_2$$

$$\begin{bmatrix} \text{Rate entropy} \\ \text{trans fer into} \\ \text{CV by heat and} \\ \text{irreversibilities} \end{bmatrix} = \begin{bmatrix} \text{Rate} \\ \text{entropy} \\ \text{leaves} \\ \text{CV} \end{bmatrix} - \begin{bmatrix} \text{Rate} \\ \text{entropy} \\ \text{enters} \\ \text{CV} \end{bmatrix}$$

$$\int_{1}^{2} \frac{\delta \dot{Q}_{\rm in}}{T} + \dot{S}_{\rm gen}$$





reversible process 1-2



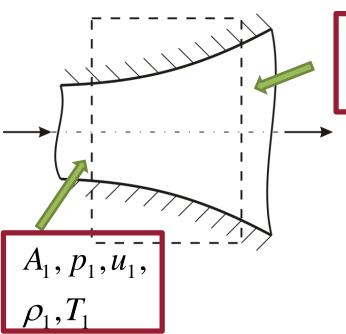
$$\int_{1}^{2} \frac{\delta Q_{\text{in}}}{T} + \dot{S}_{\text{gen}} = \dot{S}_{2} - \dot{S}_{1}$$



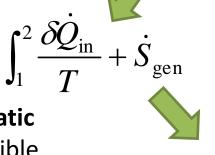


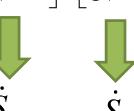
Fundamental equations

Entropy equation



$$\begin{bmatrix} \text{Rate entropy} \\ \text{trans fer into} \\ \text{CV by heat and} \\ \text{irreversibilities} \end{bmatrix} = \begin{bmatrix} \text{Rate} \\ \text{entropy} \\ \text{leaves} \\ \text{CV} \end{bmatrix} - \begin{bmatrix} \text{Rate} \\ \text{entropy} \\ \text{enters} \\ \text{CV} \end{bmatrix}$$







adiabatic

reversible

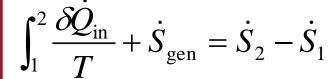
process 1-2



isentropic

process 1-2



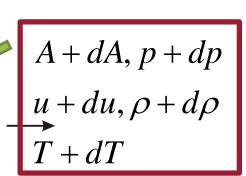




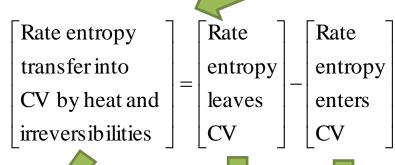


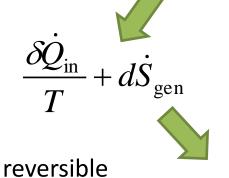
Fundamental equations

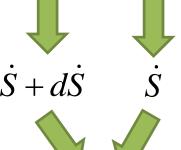
Entropy equation



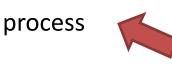
inviscid steady-state flow

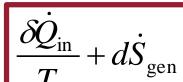






 $\frac{\delta \dot{Q}_{\rm in}}{T} = d\dot{S}, \frac{\delta \dot{q}_{\rm in}}{T} = d\dot{S}$









Fundamental equations

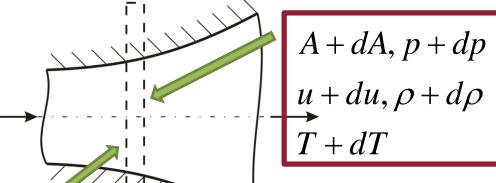
inviscid steady-state flow

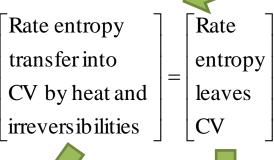
Rate

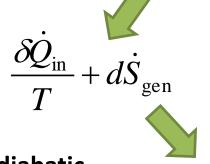
entropy

enters

Entropy equation













isentropic process

$$0 = d\dot{S}, 0 = d\dot{s}, 0 = ds$$



adiabatic reversible process



$$\frac{\delta \dot{Q}_{\rm in}}{T} + d\dot{S}_{\rm gen} = d\dot{S}$$

Quasi-one dimension fluid

flow equations:

Continuity equation



Energy equation

Entropy equation



Fundamental equations

inviscid steady-state flow

$$\frac{d\rho}{\rho} + \frac{du}{u} + \frac{dA}{A} = 0$$

$$-\frac{dp}{\rho} = udu$$

$$dh + udu = 0$$

$$0 = ds$$

adiabatic flow, no shaft work

isentropic

Thermodynamics of gases

- equation of state (EOS)
$$p = \rho RT$$
 or $p = \rho R_M T / M_m$ $pv = RT$



Thermodynamics of gases

- equation of state (EOS)
$$p = \rho RT$$
 or $p = \rho R_M T / M_m$ $pv = RT$

- internal energy
$$e$$
- enthalpy $h = e + pv$

$$e = e(T)$$

$$h = h(T)$$

$$de = c_v dT$$

$$dh = c_p dT$$

$$e = e(T)$$

$$h = h(T)$$

$$de =$$

$$dh = c_p dT$$



Thermodynamics of gases

- equation of state (EOS)
$$p = \rho RT$$
 or $p = \rho R_M T / M_m$ $pv = RT$

$$-$$
 internal energy e

- enthalpy
$$h = e + pv$$

$$e = e(T)$$
 $h = h(T)$
 $e = c_v T$
 $h = c_p T$
 $de = c_v dT$
 $dh = c_p dT$
 $calorically$
 $dh = c_p T$





Thermodynamics of gases

- perfect gas
 - $p = \rho RT$ or $p = \rho R_M T / M_m$ pv = RT $pv = R_M T / M_m$ – equation of state (EOS)
 - internal energy e
 - enthalpy h = e + pv



$$e = e(T)$$
 $h = h(T)$

$$h = h(T)$$

$$e = c_v T$$

$$e = c_{v}T$$
$$h = c_{p}T$$

$$de = c_v dT$$

$$de = c_{v}dT$$
$$dh = c_{p}dT$$

- heat capacity ratio $\gamma = c_n / c_v$
- difference between heat capacities

$$R = c_p - c_v$$





Thermodynamics of gases

- equation of state (EOS)
$$p = \rho RT$$
 or $p = \rho R_M T / M_m$ $pv = RT$

- internal energy e
- enthalpy h = e + pv



$$e = e(T)$$

$$h = h(T)$$

$$e = c_v T$$

$$h = c_p T$$

$$de = c_v dT$$

$$de = c_{v}dT$$
$$dh = c_{p}dT$$

$$\gamma = c_p / c_v$$

$$R = c_p - c_v$$

$$c_p = \frac{\gamma R}{\gamma - 1} \quad c_v =$$





Thermodynamics of gases

- moving fluid with speed u
 - static parameters



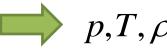


 p,T,ρ "measured" in moving fluid



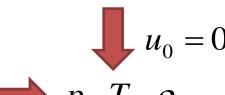
Thermodynamics of gases

- moving fluid with speed u
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 p,T,ρ "measured" in moving fluid

fluid brought to rest adiabatically







Thermodynamics of gases

- moving fluid with speed u
 - static parameters

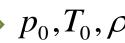




fluid brought to rest adiabatically



$$u_0 = 0$$





total (stagnation) pressure, temperature, density

 corresponding total enthalpy to total temperature



$$h_0 = c_p T_0$$

Thermodynamics of gases

- stagnation conditions
 - adiabatic process

$$\left(h_1 + \frac{u_1^2}{2}\right) = \text{const.} \qquad \text{rest of fluid} \qquad h_0 = \left(h_1 + \frac{u_1^2}{2}\right)$$



total enthalpy is constant through steady, inviscid, adiabatic flow

$$h_0 = \left(h_1 + \frac{u_1^2}{2}\right)$$

Thermodynamics of gases

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$$h_0 = \left(h_1 + \frac{u_1^2}{2}\right)$$

also total temperature



$$h_0 = c_p T_0$$



Thermodynamics of gases

- stagnation conditions
 - adiabatic process

total enthalpy is constant through steady, inviscid, adiabatic flow

$$\left(h_1 + \frac{u_1^2}{2}\right) = \text{const.}$$
 rest of fluid

$$h_0 = \left(h_1 + \frac{u_1^2}{2}\right)$$

also total temperature is constant

$$h_0 = c_p T$$

isentropic process



$$p_0, \rho_0$$



total pressure and total density are constant



Thermodynamics of gases

- moving fluid with speed u
 - static parameters





 p,T,ρ "measured" in moving fluid



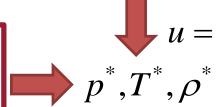
Thermodynamics of gases

- moving fluid with speed u
 - static parameters



moving fluid

fluid brought to speed of sound isentropically



Thermodynamics of gases

- moving fluid with speed u
 - static parameters



moving fluid

fluid brought to speed of sound isentropically



$$p^*,T^*,\rho^*$$



critical pressure, temperature, density



Thermodynamics of gases

- reversible process
 - added heat closed system \longrightarrow $\delta w_{\rm in} + \delta q_{\rm in} = de$

 - internal energy $de = c_v dT$
 - pressure work $\delta w_{\rm in} = -pdv$



Thermodynamics of gases

using

reversible process

- added heat - closed system
$$\longrightarrow$$
 $\delta w_{\rm in} + \delta q_{\rm in} = de$

- reversible process
$$\delta q_{\rm in} = dsT$$

- internal energy $de = c_v dT$

— internal energy
$$\implies de = c_v dT$$

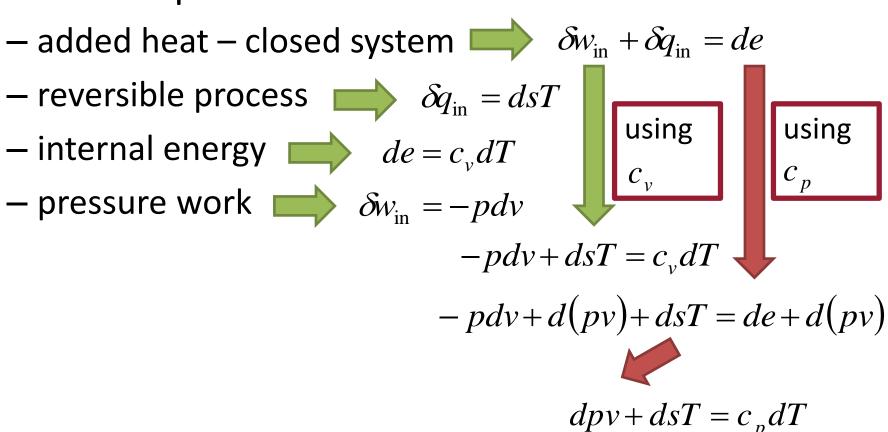
- pressure work
$$\longrightarrow \delta w_{\rm in} = -pdv$$

$$-pdv + dsT = c_v dT$$



Thermodynamics of gases

reversible process



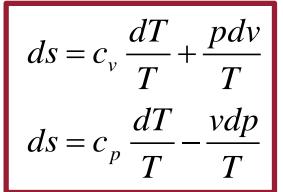
Thermodynamics of gases

- reversible process
 - added heat closed system \longrightarrow $\delta w_{in} + \delta q_{in} = de$



- internal energy $de = c_v dT$
- pressure work $\longrightarrow \delta w_{in} = -pdv$

$$\delta w_{\rm in} = -pdv$$



using

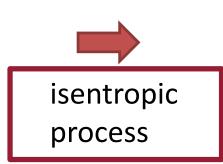
Thermodynamics of gases

- reversible process
 - added heat closed system $\longrightarrow \delta w_{\rm in} + \delta q_{\rm in} = de$

 - internal energy $de = c_v dT$
 - pressure work $\longrightarrow \delta w_{\rm in} = -pdv$

$$ds = c_v \frac{dT}{T} + \frac{pdv}{T}$$

$$ds = c_p \frac{dT}{T} - \frac{vdp}{T}$$



$$0 = c_v \frac{dI}{T} + \frac{pdv}{T}$$
$$0 = c_p \frac{dT}{T} - \frac{vdp}{T}$$



Thermodynamics of gases

- reversible process
 - added heat closed system \longrightarrow $\delta w_{\rm in} + \delta q_{\rm in} = de$

 - internal energy $de = c_v dT$
 - pressure work $\longrightarrow \delta w_{\rm in} = -pdv$

$$\frac{v_2}{v_1} = \left(\frac{T_2}{T_1}\right)^{\frac{-1}{\gamma-1}}$$

$$0 = c_v \frac{dT}{T} + \frac{Rdv}{v}$$

$$0 = c_v \frac{dT}{T} + \frac{pdv}{T}$$

$$\frac{p_2}{p_1} = \left(\frac{T_2}{T_1}\right)^{\frac{\gamma}{\gamma-1}}$$

$$0 = c_v \frac{dT}{T} - \frac{Rdp}{T}$$

$$0 = c_v \frac{dT}{T} - \frac{vdp}{T}$$

$$0 = c_v \frac{dT}{T} - \frac{vdp}{T}$$

$$v = c_p = \frac{\gamma R}{\gamma - 1}$$





Thermodynamics of gases

- reversible process
 - added heat closed system \longrightarrow $\delta w_{in} + \delta q_{in} = de$
 - reversible process $\delta q_{\rm in} = dsT$
 - internal energy $de = c_v dT$
 - pressure work $\longrightarrow \delta w_{\rm in} = -pdv$

$$\frac{v_2}{v_1} = \left(\frac{T_2}{T_1}\right)^{\frac{-1}{\gamma - 1}}$$

$$\frac{p_2}{p_1} = \left(\frac{T_2}{T_1}\right)^{\frac{\gamma}{\gamma - 1}}$$

$$\frac{v_2}{v_1} = \left(\frac{T_2}{T_1}\right)^{\frac{-1}{\gamma - 1}}$$

$$\frac{p_2}{p_1} = \left(\frac{T_2}{T_1}\right)^{\frac{\gamma}{\gamma - 1}} = \left(\frac{v_1}{v_2}\right)^{\gamma} = \left(\frac{\rho_2}{\rho_1}\right)^{\gamma}$$

$$\frac{p_2}{p_1} = \left(\frac{T_2}{T_1}\right)^{\frac{\gamma}{\gamma - 1}} = \left(\frac{v_1}{v_2}\right)^{\gamma} = \left(\frac{\rho_2}{\rho_1}\right)^{\gamma}$$

isentropic process



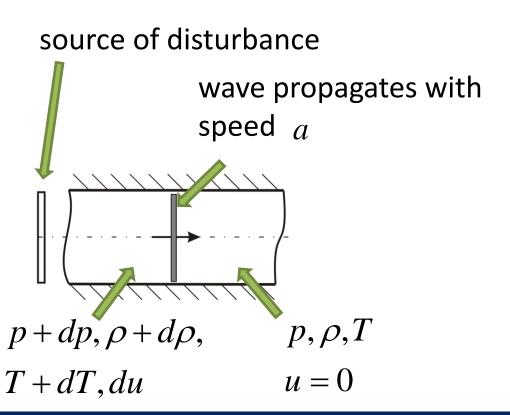


Speed of sound

- physical mechanism:
 - sound propagation in gas is based on molecular motion
 - energy is transfer to gas molecules, they start to move about in random fashion
 - they collide with other molecules and transfer their energy to these molecules
 - the process of collision repeats energy is propagated
 - macroscopic parameters p,T,ρ are slightly varied by increased microscopic parameter energy of molecule



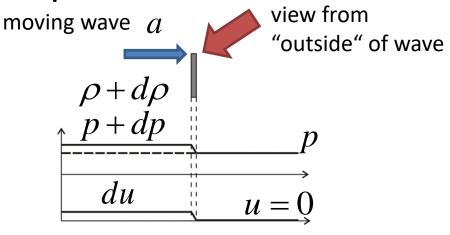
Speed of sound

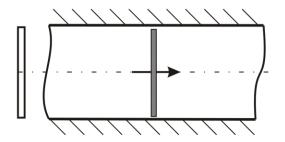




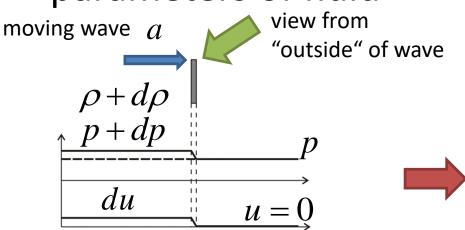


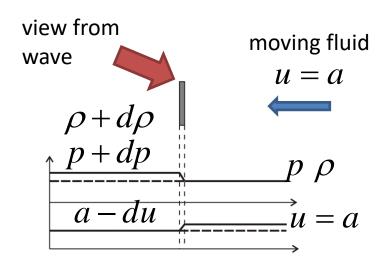
Speed of sound

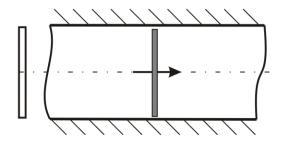




Speed of sound



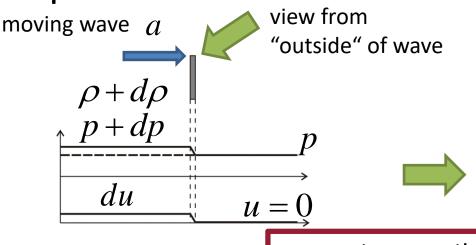






Speed of sound

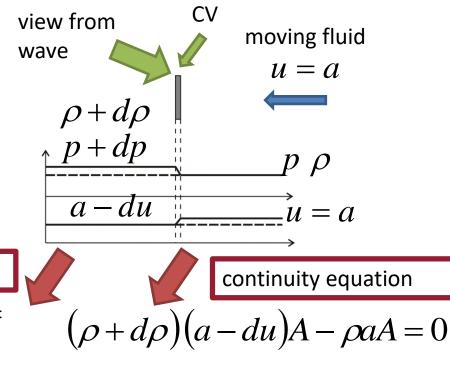
parameters of fluid



momentum equation

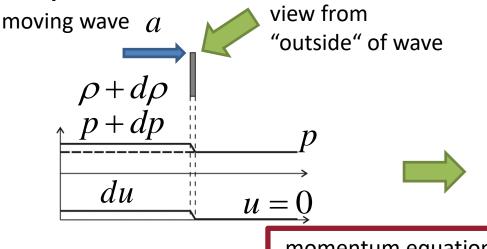
$$(\rho + d\rho)(a - du) A - \rho a^2 A =$$

$$= A(p - (p + dp))$$



Speed of sound

parameters of fluid

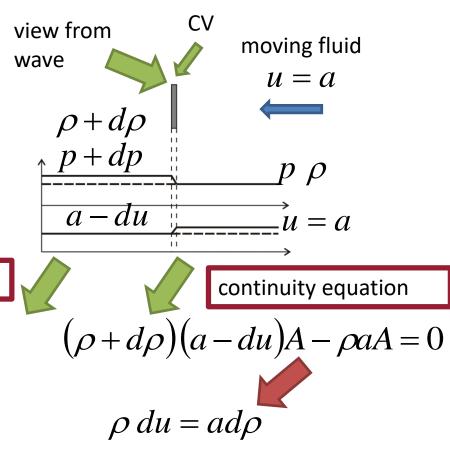


momentum equation

$$(\rho + d\rho)(a - du) A - \rho a^{2} A =$$

$$= A(p - (p + dp))$$

$$a = dp$$



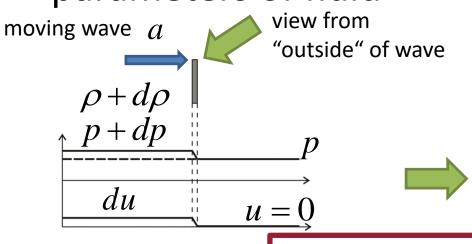




Speed of sound

continuity equation

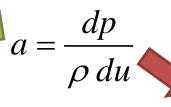
parameters of fluid



momentum equation $(\rho + d\rho)(a - du) A - \rho a^2 A =$ = A(p - (p + dp))

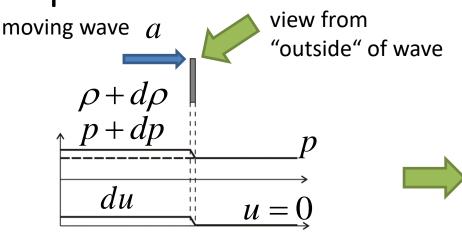
$$(\rho + d\rho)(a - du)A - \rho aA = 0$$

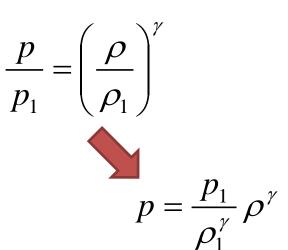
$$\rho du = ad\rho$$

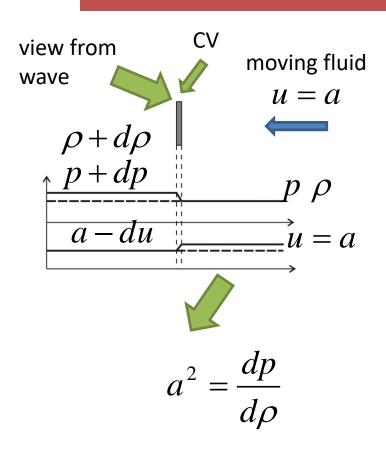


$$a^2 = \frac{dp}{d\rho}$$

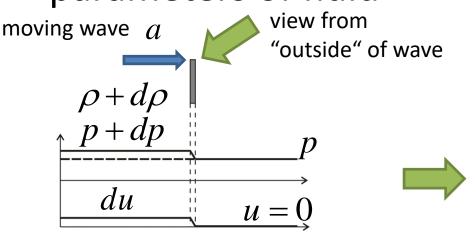
Speed of sound



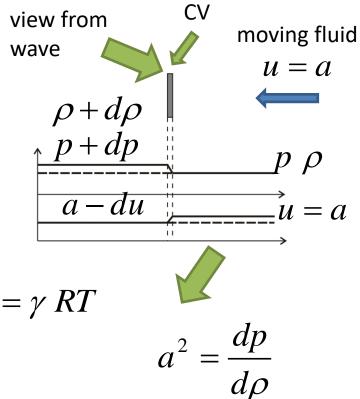




Speed of sound



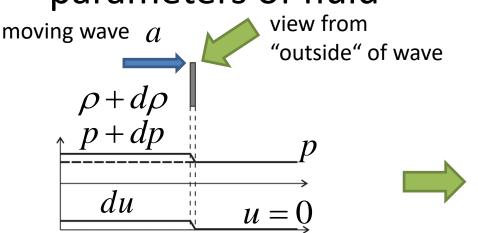
$$\frac{p}{p_1} = \left(\frac{\rho}{\rho_1}\right)^{\gamma} \quad \frac{dp}{d\rho} = \frac{p_1}{\rho_1^{\gamma}} \gamma \rho^{\gamma - 1} = \frac{p}{\rho} \gamma = \gamma RT$$

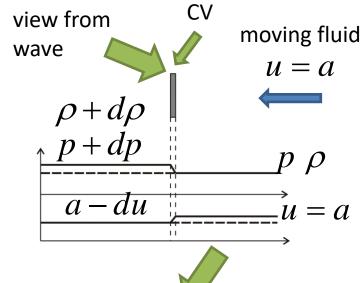


$$p = \frac{p_1}{\rho_1^{\gamma}} \rho^{\gamma}$$

Speed of sound

parameters of fluid





$$\frac{p}{p_1} = \left(\frac{\rho}{\rho_1}\right)^{\gamma} \quad \frac{dp}{d\rho} = \frac{p_1}{\rho_1^{\gamma}} \gamma \rho^{\gamma - 1} = \frac{p}{\rho} \gamma = \gamma RT$$

$$p_1$$

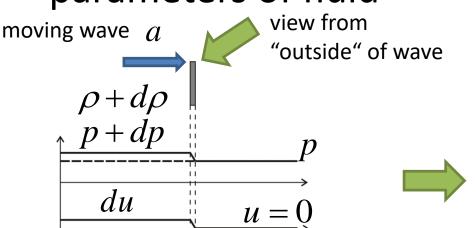
$$a = \sqrt{\gamma RT}$$

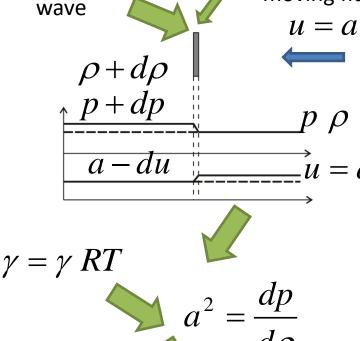


Speed of sound

moving fluid

parameters of fluid





view from

$$\frac{p}{p_1} = \left(\frac{\rho}{\rho_1}\right)^{\gamma} \quad \frac{dp}{d\rho} = \frac{p_1}{\rho_1^{\gamma}} \gamma \rho$$

$$p_1 \quad \gamma$$

$$a = \sqrt{\gamma RT}$$

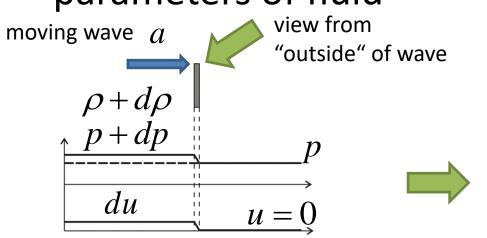
$$M = \frac{u}{a} = \frac{u}{\sqrt{\gamma RT}}$$



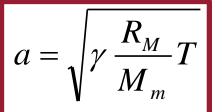


Speed of sound

parameters of fluid

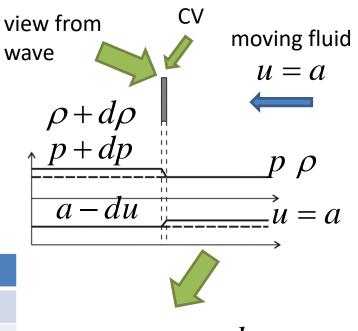


Gas	M _m [g/mol]	γ [–]	a at 0°C [m/s]
Air	28.96	1.404	331
Hydrogen	2.016	1.407	1270
Xenon	131.3	1.667	170





$$a = \sqrt{\gamma RT}$$



$$a^2 = \frac{dp}{d\rho}$$

$$M = \frac{u}{a} = \frac{u}{\sqrt{\gamma RT}}$$





Isentropic flow

- Importance of isentropic flow:
 - isentropic flow is adiabatic in which viscous losses are negligible
 - real flows are not isentropic



the effects of viscosity and heat transfer are restricted to thin layers near the walls

major part of the flow can be assumed to be isentropic



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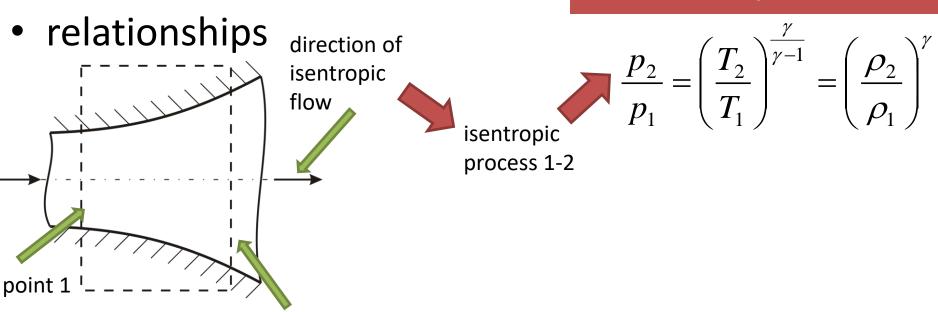


many flows in engineering practice can be adequately modeled by assuming them to be **isentropic** and also **steady-state** and **quasi-one dimensional flow**





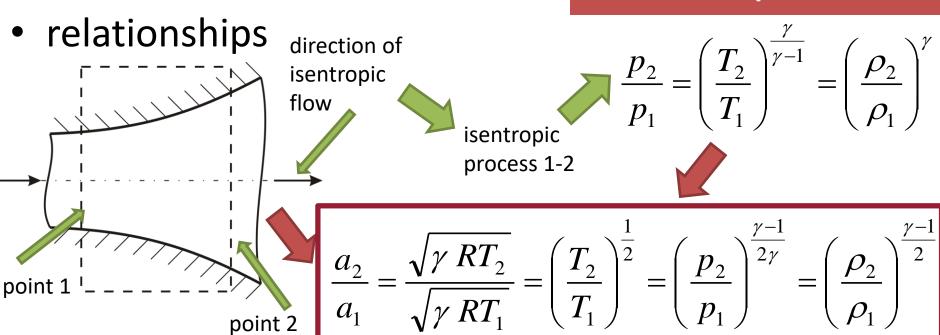
Isentropic flow





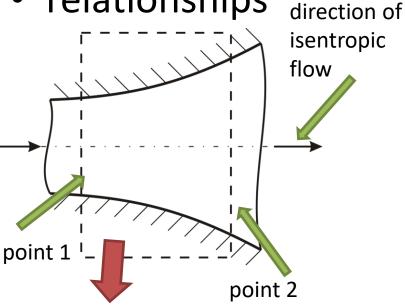
point 2

Isentropic flow



Isentropic flow





adiabatic energy equation 1-2

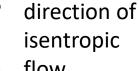
$$\left(h_1 + \frac{u_1^2}{2}\right) = \left(h_2 + \frac{u_2^2}{2}\right)$$

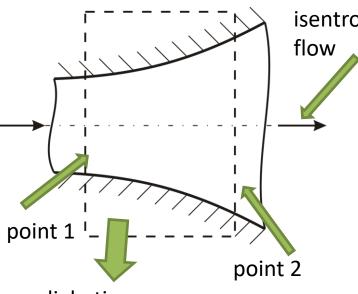




Isentropic flow







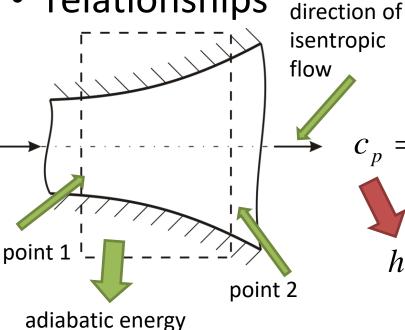
$$h = c_p T$$

adiabatic energy equation 1-2

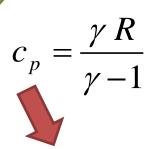
$$\left(h_1 + \frac{u_1^2}{2}\right) = \left(h_2 + \frac{u_2^2}{2}\right)$$

Isentropic flow





$$\left(h_1 + \frac{u_1^2}{2}\right) = \left(h_2 + \frac{u_2^2}{2}\right)$$

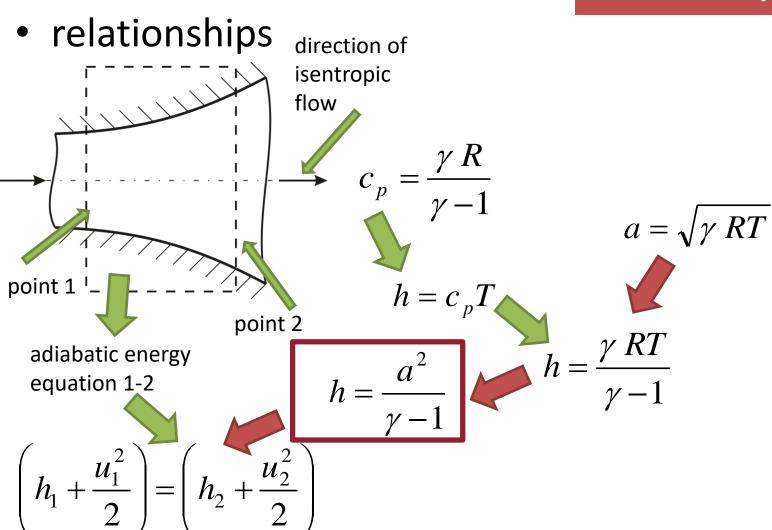


$$h = c_p T$$

$$h = \frac{\gamma RT}{\gamma - 1}$$

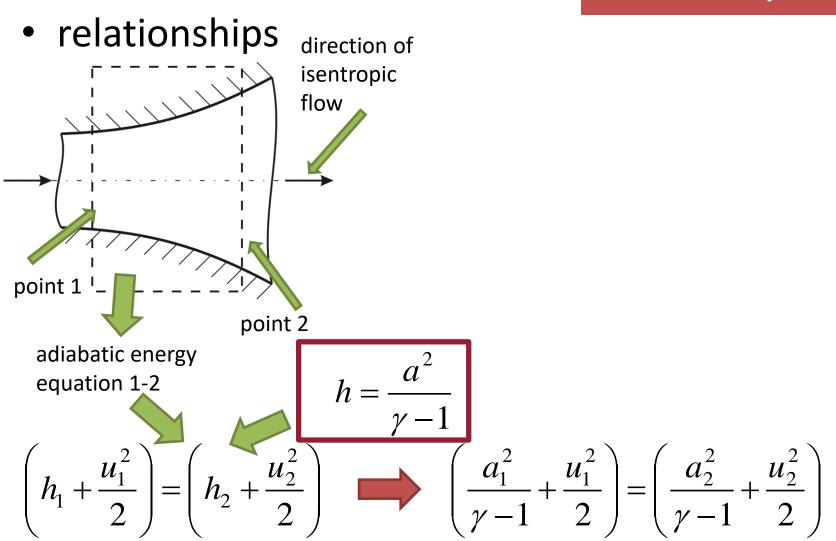
equation 1-2

Isentropic flow





Isentropic flow

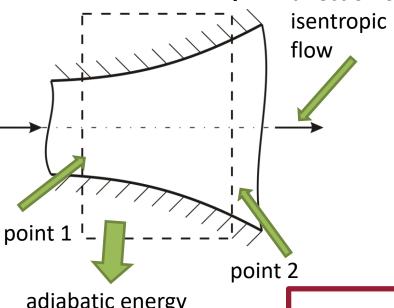






Isentropic flow





adiabatic energy equation 1-2

$$h = \frac{a^2}{\gamma - 1}$$

$$\left(h_1 + \frac{u_1^2}{2}\right) = \left(h_2 + \frac{u_2^2}{2}\right)$$



$$\left(\frac{a^2}{u^2(\gamma - 1)} + \frac{1}{2}\right) = \frac{\gamma + 1}{2(\gamma - 1)} \frac{a^{*2}}{u^2}$$



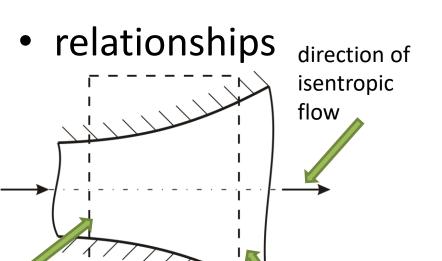
$$\left(\frac{a^2}{\gamma - 1} + \frac{u^2}{2}\right) = \frac{\gamma + 1}{2(\gamma - 1)}a^{*2}$$





$$\left(\frac{a_1^2}{\gamma - 1} + \frac{u_1^2}{2}\right) = \left(\frac{a_2^2}{\gamma - 1} + \frac{u_2^2}{2}\right)$$

Isentropic flow



$$\left(\frac{a^{2}}{u^{2}(\gamma-1)} + \frac{1}{2}\right) = \frac{\gamma+1}{2(\gamma-1)} \frac{a^{*2}}{u^{2}}$$

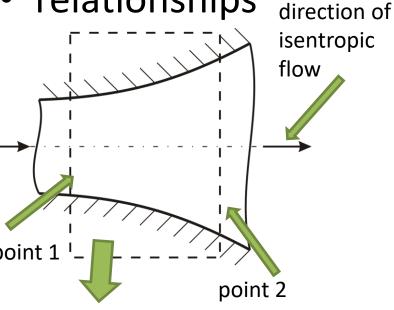
$$\left(\frac{1}{M^{2}(\gamma-1)}+\frac{1}{2}\right)=\frac{\gamma+1}{2(\gamma-1)}\frac{1}{\left(M^{*}\right)^{2}}$$

point 1

point 2

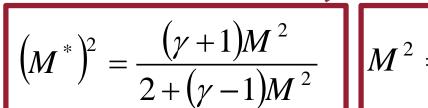
Isentropic flow

relationships



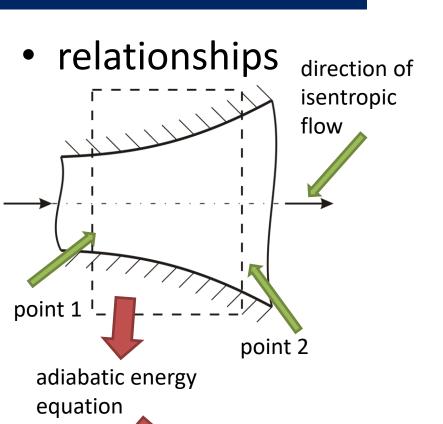
$$\left(\frac{a^2}{u^2(\gamma - 1)} + \frac{1}{2}\right) = \frac{\gamma + 1}{2(\gamma - 1)} \frac{a^{*2}}{u^2}$$

$$\left(\frac{1}{M^{2}(\gamma-1)} + \frac{1}{2}\right) = \frac{\gamma+1}{2(\gamma-1)} \frac{1}{(M^{*})^{2}}$$



$$M^{2} = \frac{2}{(\gamma + 1)/(M^{*})^{2} - (\gamma - 1)}$$

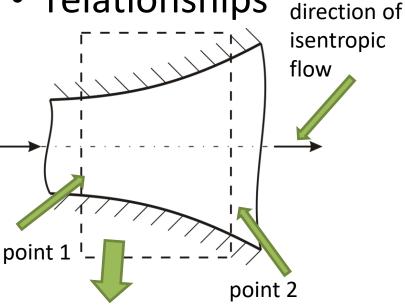
Isentropic flow



$$\left(h + \frac{u^2}{2}\right) = h_0 \qquad \left(c_p T + \frac{u^2}{2}\right) = c_p T_0$$

Isentropic flow





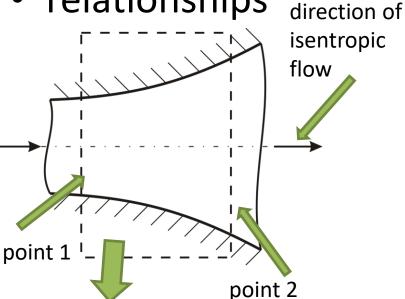
adiabatic energy equation

$$\left(h + \frac{u^2}{2}\right) = h_0 \qquad \left(c_p T + \frac{u^2}{2}\right) = c_p T_0$$

$$1 + \frac{u^2}{2c_p T} = \frac{T_0}{T}$$

Isentropic flow





$$1 + \frac{(\gamma - 1)u^2}{2\gamma RT} = \frac{T_0}{T}$$

adiabatic energy

$$\left(h + \frac{u^2}{2}\right) = h_0 \Longrightarrow \left(c_p T + \frac{u^2}{2}\right) = c_p T_0$$

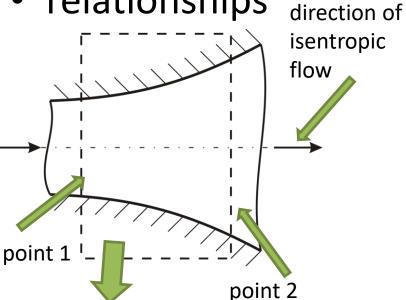
$$\left(c_p T + \frac{u^2}{2}\right) = c_p T_0$$

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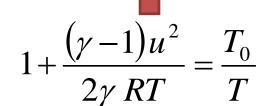
$$c_p = \frac{\gamma R}{\gamma - 1}$$

Isentropic flow





$$1 + \frac{\left(\gamma - 1\right)}{2}M^2 = \frac{T_0}{T}$$



adiabatic energy

equation

$$\left(h + \frac{u^2}{2}\right) = h_0 \qquad \left(c_p T + \frac{u^2}{2}\right) = c_p T_0$$

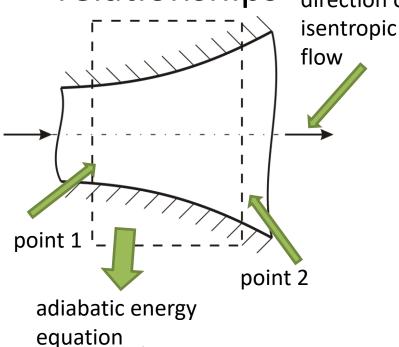
$$1 + \frac{u^2}{2c_p T} = \frac{T_0}{T}$$

$$c_p = \frac{\gamma R}{\gamma - 1}$$



Isentropic flow

relationships direction of



$$\left(1 + \frac{(\gamma - 1)}{2} M^{2}\right)^{\frac{1}{\gamma - 1}} = \frac{\rho_{0}}{\rho}$$

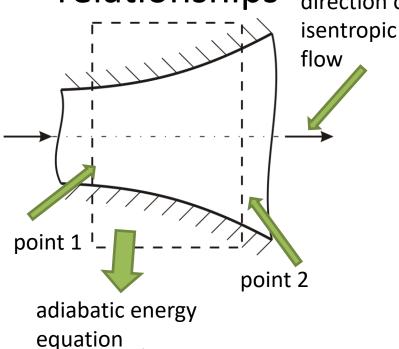
$$\left(1 + \frac{(\gamma - 1)}{2} M^{2}\right)^{\frac{\gamma}{\gamma - 1}} = \frac{p_{0}}{p}$$

$$1 + \frac{(\gamma - 1)}{2} M^{2} = \frac{T_{0}}{T}$$

$$\left(h + \frac{u^2}{2}\right) = h_0 \qquad \left(c_p T + \frac{u^2}{2}\right) = c_p T_0 \qquad \frac{p_0}{p} = \left(\frac{T_0}{T}\right)^{\frac{\gamma}{\gamma - 1}} = \left(\frac{\rho_0}{\rho}\right)^{\gamma}$$

Isentropic flow

relationships direction of



$$\left(1 + \frac{(\gamma - 1)}{2}M^{2}\right)^{\frac{1}{\gamma - 1}} = \frac{\rho_{0}}{\rho}$$

$$\left(1 + \frac{(\gamma - 1)}{2}M^{2}\right)^{\frac{\gamma}{\gamma - 1}} = \frac{p_{0}}{p}$$

$$1 + \frac{(\gamma - 1)}{2}M^{2} = \frac{T_{0}}{T}$$

if
$$M = 1$$

$$\rho = \rho^*$$

$$p = p^*$$

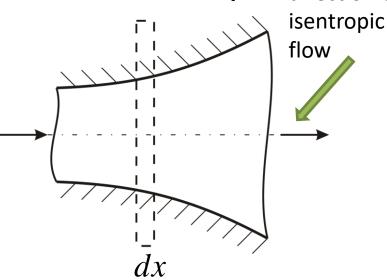
$$T = T^*$$

$$\left(h + \frac{u^2}{2}\right) = h_0 \qquad \left(c_p T + \frac{u^2}{2}\right) = c_p T_0 \qquad \frac{p_0}{p} = \left(\frac{T_0}{T}\right)^{\frac{\gamma}{\gamma - 1}} = \left(\frac{\rho_0}{\rho}\right)^{\frac{\gamma}{\gamma}}$$

direction of

Isentropic flow

relationships

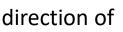


momentum equation

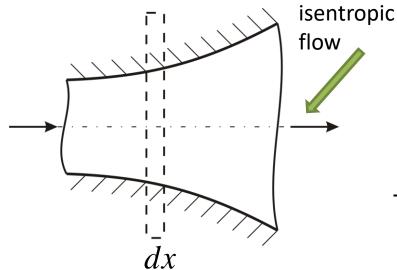
$$-\frac{dp}{\rho} = udu$$

Isentropic flow

relationships



momentum equation



$$-\frac{dp}{\rho} = udu$$

$$-\frac{dp}{p} = \frac{\rho u}{pu} u du$$



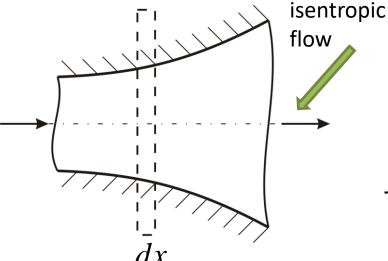
Isentropic flow

relationships

direction of

momentum equation

speed of sound



$$-\frac{dp}{\rho} = udu \qquad a^2 = \gamma RT = \gamma \frac{p}{\rho}$$

$$-\frac{dp}{p} = \frac{\rho u}{pu}udu$$

$$a^2 = \gamma RT = \gamma \frac{p}{\rho}$$



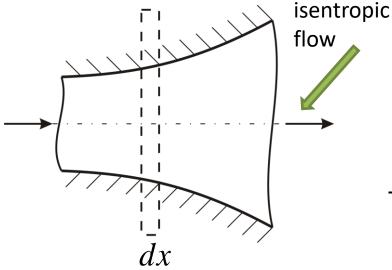
Isentropic flow

relationships

direction of

momentum equation

speed of sound



$$-\frac{dp}{\rho} = udu \qquad a^2 = \gamma RT = \gamma \frac{p}{\rho}$$

$$-\frac{dp}{p} = \frac{\rho u}{pu} u du$$



$$\frac{dp}{p} = -\gamma \frac{u^2}{a^2} \frac{du}{u}$$





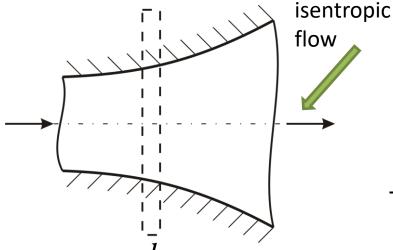
Isentropic flow

relationships

direction of

momentum equation

speed of sound



$$-\frac{dp}{\rho} = udu$$

$$a^2 = \gamma RT = \gamma \frac{p}{\rho}$$

$$-\frac{dp}{p} = \frac{\rho u}{pu} u du$$



 magnitude of fractional pressure change induced by a given fractional velocity change depends on square of Mach number

$$\frac{dp}{p} = -\gamma \frac{u^2}{a^2} \frac{du}{u}$$

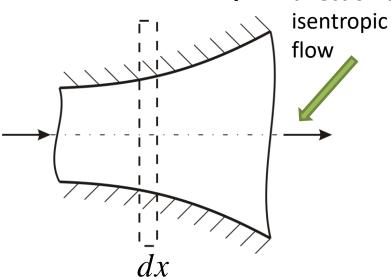


$$\frac{dp}{p} = -\gamma M^2 \frac{du}{u}$$

direction of

Isentropic flow

relationships



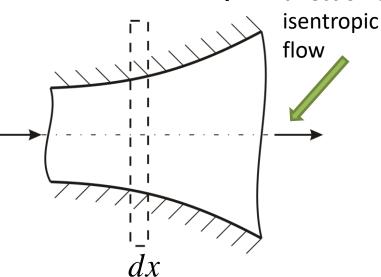
energy equation

$$dh + udu = 0$$

direction of

Isentropic flow

relationships



energy equation

$$dh + udu = 0$$



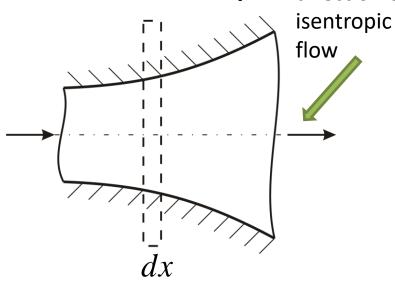
calorically perfect gas

$$dh = c_p dT$$

direction of

Isentropic flow

relationships



energy equation

$$dh + udu = 0$$

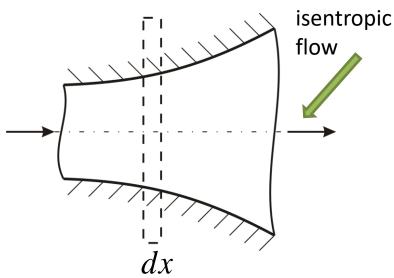
$$c_n dT = -u dv$$

calorically perfect gas $dh = c_{p}dT$



Isentropic flow

relationships



direction of energy equation

$$dh + udu = 0$$

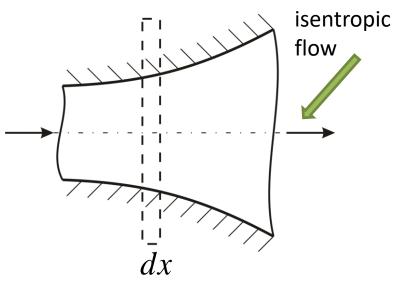
$$c_p dT = -udu$$

$$\frac{dT}{T} = -\frac{u}{c_p T} \frac{u}{u} du$$

calorically perfect gas $dh = c_{p}dT$

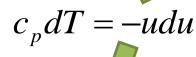
Isentropic flow

relationships



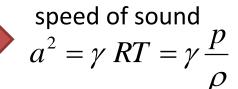
direction of energy equation

$$dh + udu = 0$$



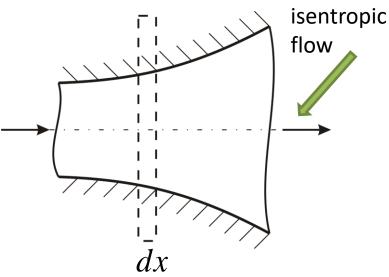
$$\frac{dT}{T} = -\frac{u}{c_{\cdot \cdot}T} \frac{u}{u} du$$

calorically perfect gas $dh = c_n dT$



Isentropic flow

relationships



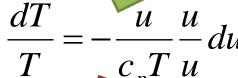
direction of

energy equation

$$dh + udu = 0$$

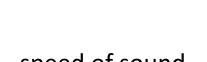


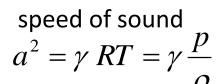




$$\frac{dT}{T} = -\frac{\gamma R u^2}{c_n a^2} \frac{du}{u}$$

calorically perfect gas $dh = c_n dT$



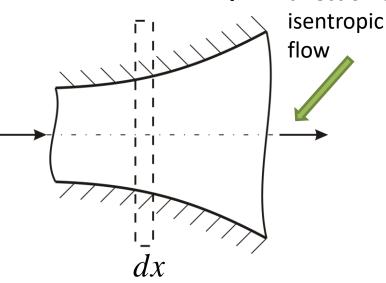






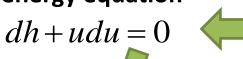
Isentropic flow

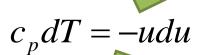
relationships



direction of

energy equation





$$\frac{dT}{T} = -\frac{u}{c \cdot T} \frac{u}{u} dt$$

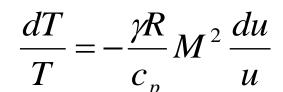
$$\frac{dT}{T} = -\frac{\gamma R u^2}{c_p a^2} \frac{du}{u}$$



calorically perfect gas $dh = c_n dT$



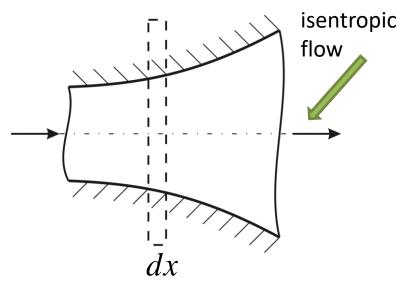
speed of sound
$$a^2 = \gamma RT = \gamma \frac{p}{\rho}$$



direction of

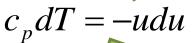
Isentropic flow

relationships



 magnitude of fractional temperature change induced by a given fractional velocity change depends on square of Mach number energy equation

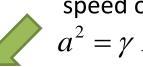
$$dh + udu = 0$$



$$\frac{dT}{T} = -\frac{u}{c_n T} \frac{u}{u} du$$

$$\frac{dT}{T} = -\frac{\gamma R u^2}{c_n a^2} \frac{du}{u}$$

calorically perfect gas $dh = c_n dT$



speed of sound
$$a^2 = \gamma RT = \gamma \frac{p}{r}$$

$$\frac{dT}{T} = -(\gamma - 1)M^2 \frac{du}{u}$$

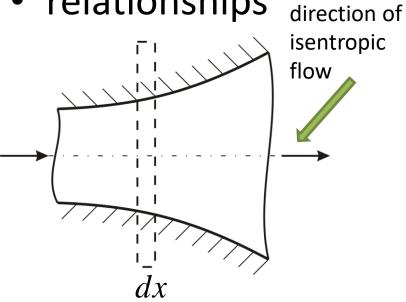
$$\frac{dT}{T} = -\frac{\gamma R}{c_n} M^2 \frac{du}{u}$$





Isentropic flow

relationships

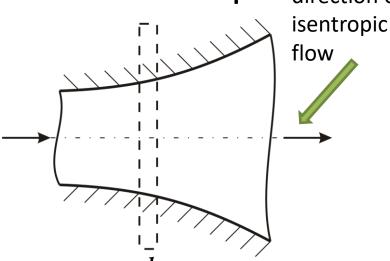


equation of state

$$\frac{p}{\rho T} = R$$

Isentropic flow

relationships



direction of equation of state

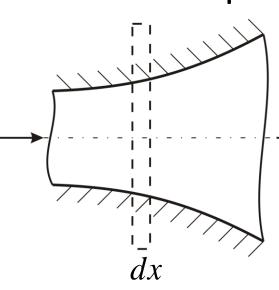
$$\frac{dp}{p} - \frac{d\rho}{\rho} - \frac{dT}{T} = 0$$

$$\frac{dp}{p} - \frac{d\rho}{\rho} - \frac{dT}{T} = 0$$

$$\frac{dp}{p} = -\gamma M^2 \frac{du}{u}$$

Isentropic flow

relationships



direction of equation of state

isentropic flow
$$\frac{p}{\rho T} = R \qquad \frac{dT}{T} = -(\gamma - 1)M^{2} \frac{d}{u}$$

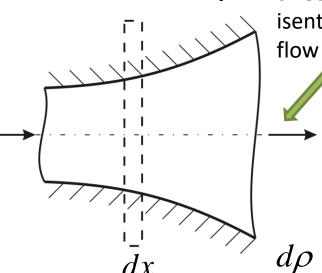
$$\frac{dp}{p} - \frac{d\rho}{\rho} - \frac{dT}{T} = 0$$

$$\frac{dp}{p} = -\gamma M^{2} \frac{du}{u}$$

$$\frac{d\rho}{\rho} = -\gamma M^{2} \frac{du}{u} + (\gamma - 1)M^{2} \frac{du}{u}$$

Isentropic flow

relationships



direction of

isentropic

equation of state

$$\frac{dp}{\rho T} = R$$

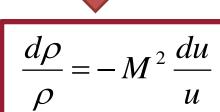
$$\frac{dT}{T} = -(\gamma - 1)M^{2}$$

$$\frac{dp}{p} - \frac{d\rho}{\rho} - \frac{dT}{T} = 0$$

$$\frac{dp}{p} = -\gamma M^{2} \frac{du}{u}$$

$$\frac{d\rho}{\rho} = -\gamma M^2 \frac{du}{u} + (\gamma - 1)M^2 \frac{du}{u}$$

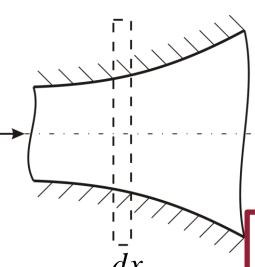
 magnitude of fractional temperature change induced by a given fractional velocity change depends on square of Mach number



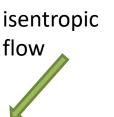


Isentropic flow

relationships



direction of



 magnitude of fractional properties change induced by a given fractional velocity change depends on square of Mach number

$$\frac{dp}{p} = -\gamma M^2 \frac{du}{u}$$

$$\frac{dT}{T} = -(\gamma - 1)M^2 \frac{du}{u}$$

$$\frac{d\rho}{dt} = -M^2 \frac{du}{dt}$$



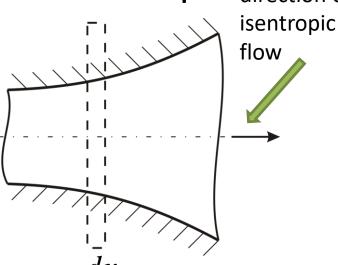
$$\frac{dp/p}{du/u} = -\gamma M^{2}$$

$$\frac{dT/T}{du/u} = -(\gamma - 1)M^{2}$$

$$\frac{d\rho/\rho}{du/u} = -M^{2}$$

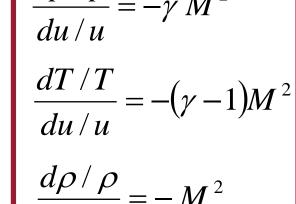
Isentropic flow

relationships direction of



• magnitude of fractional properties change induced by a given fractional velocity change depends on square of Mach number

	fractional propert. change induced by fractional velocity change of air [%]		
Mach num.	density	temp.	pressure
0.1	1	1.4	0.4
0.33	10.9	15.2	4.4
0.4	16	22.4	6.4







Nozzle fluid flow

governing equations for analysis of nozzle

continuity

momentum

energy

state

$$\frac{d\rho}{\rho} + \frac{du}{u} + \frac{dA}{A} = 0$$

$$-\frac{dp}{\rho} = udu$$

$$dh + udu = 0$$

$$\frac{dp}{\rho} - \frac{d\rho}{\rho} - \frac{dT}{T} = 0$$

Nozzle fluid flow

$$\frac{d\rho}{\rho} + \frac{du}{u} + \frac{dA}{A} = 0$$

$$\frac{d\rho}{\rho} = -M^2 \frac{du}{u} \qquad -\frac{dp}{\rho} = udu$$

$$\frac{dh + udu = 0}{\rho}$$

$$\frac{dp}{\rho} - \frac{d\rho}{\rho} - \frac{dT}{T} = 0$$



Nozzle fluid flow

$$\frac{d\rho}{\rho} + \frac{du}{u} + \frac{dA}{A} = 0 \qquad -M^2 \frac{du}{u} + \frac{du}{u} + \frac{dA}{A} = 0$$

$$\frac{d\rho}{\rho} = -M^2 \frac{du}{u} \qquad -\frac{dp}{\rho} = udu$$

$$dh + udu = 0$$

$$\frac{dp}{p} - \frac{d\rho}{\rho} - \frac{dT}{T} = 0$$



Nozzle fluid flow

$$\frac{d\rho}{\rho} + \frac{du}{u} + \frac{dA}{A} = 0$$

$$-\frac{dp}{\rho} = udu$$

$$\frac{d\rho}{\rho} - \frac{d\rho}{\rho} - \frac{dT}{T} = 0$$

$$-M^{2} \frac{du}{u} + \frac{du}{u} + \frac{dA}{A} = 0$$

$$\frac{dA}{A} = (M^{2} - 1) \frac{du}{u}$$

Nozzle fluid flow

governing equations for analysis of nozzle

$$\frac{dA}{A} = \left(M^2 - 1\right) \frac{du}{u}$$















increase in velocity +du is associated with decrease in area -dA

increase in velocity +du is associated with increase in area +dA

dA = 0 area reaches an extremum – minimum



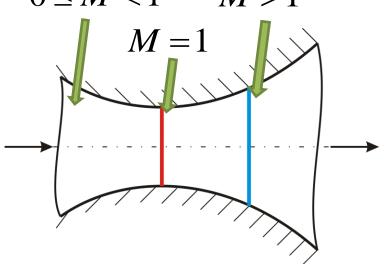
Nozzle fluid flow

governing equations for analysis of nozzle

$$\frac{dA}{A} = \left(M^2 - 1\right)\frac{du}{u}$$



 $0 \le M < 1$ M > 1



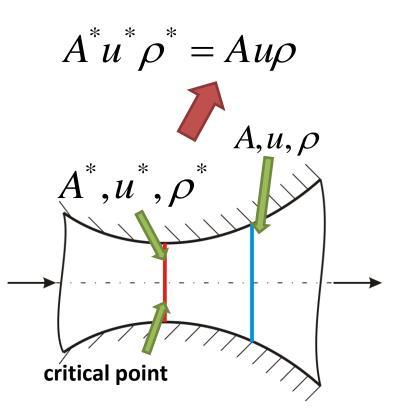
increase in velocity +du is associated with decrease in area -dA

increase in velocity +du is associated with increase in area +dA

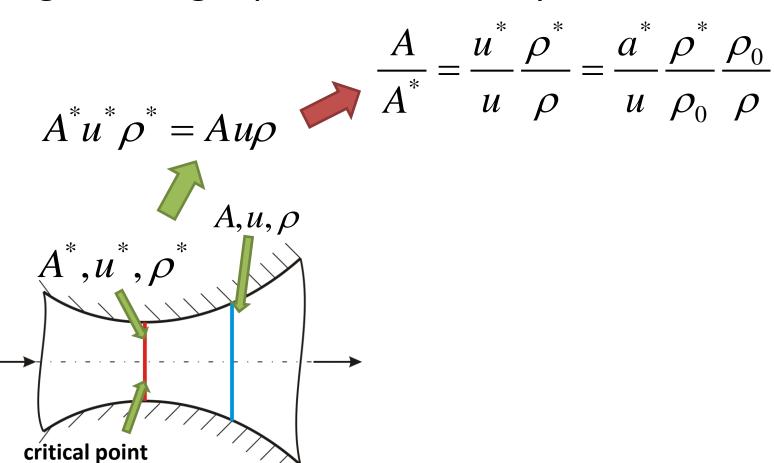
dA = 0 area reaches an extremum – minimum



Nozzle fluid flow



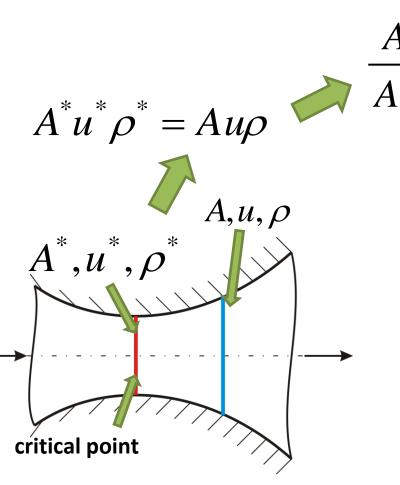
Nozzle fluid flow







Nozzle fluid flow



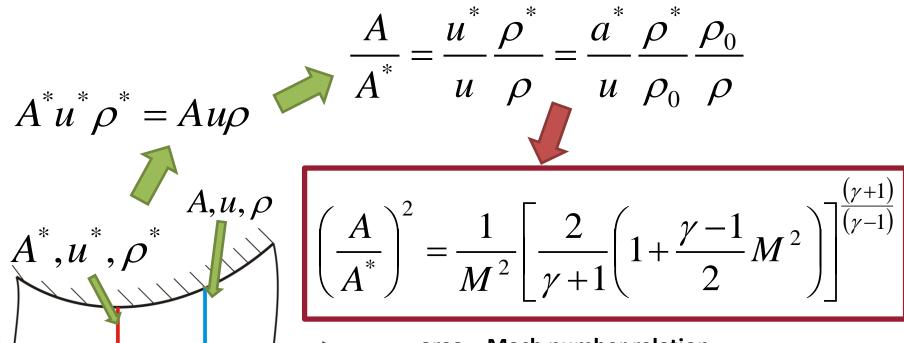
$$\frac{\rho_0}{\rho} = \left(1 + \frac{(\gamma - 1)}{2} M^2\right)^{\frac{1}{\gamma - 1}}$$

$$\frac{\rho_0}{\rho^*} = \left(1 + \frac{(\gamma - 1)}{2}\right)^{\frac{1}{\gamma - 1}}$$

$$\left(M^*\right)^2 = \left(\frac{u}{a^*}\right)^2 = \frac{(\gamma + 1)M^2}{2 + (\gamma - 1)M^2}$$

Nozzle fluid flow

governing equations for analysis of nozzle

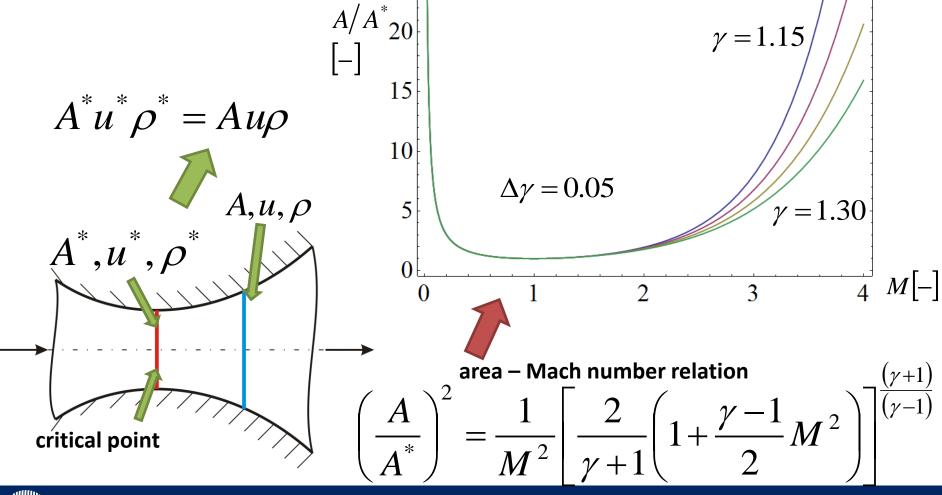


area - Mach number relation



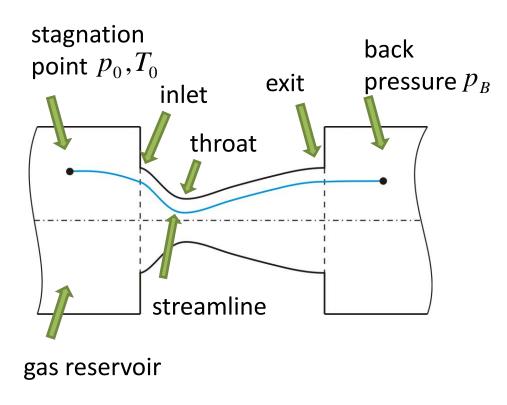
critical point

Nozzle fluid flow



Nozzle fluid flow

variation of parameters in nozzle



• stagnation parameters:

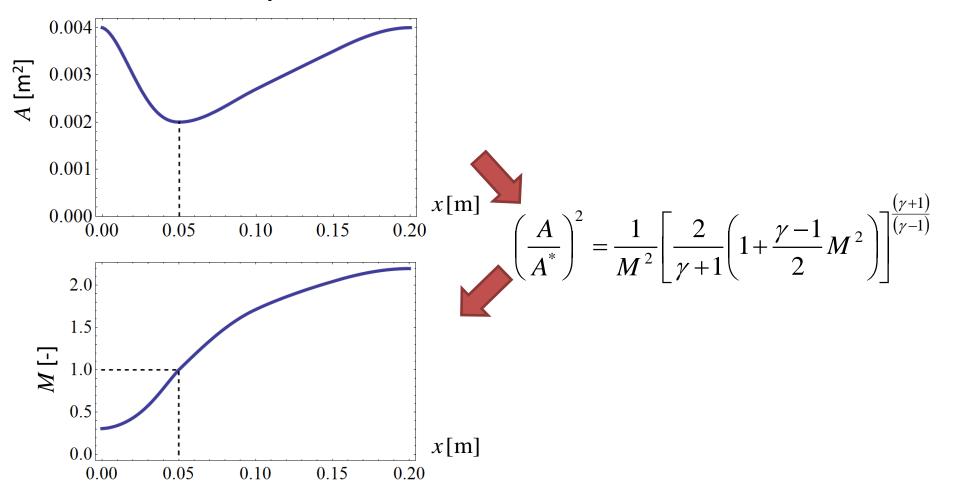
$$T_0 = 500 \,^{\circ} \text{K}$$
$$p_0 = 1 \text{MPa}$$

- nozzle area: inlet area (location 0 m): 0.004 m² throat area (loc. 0.05 m): 0.002 m²
 - exit area (loc. 0.2 m): 0.004 m²
- air: R = 288 J/kgK $\gamma = 1.4$



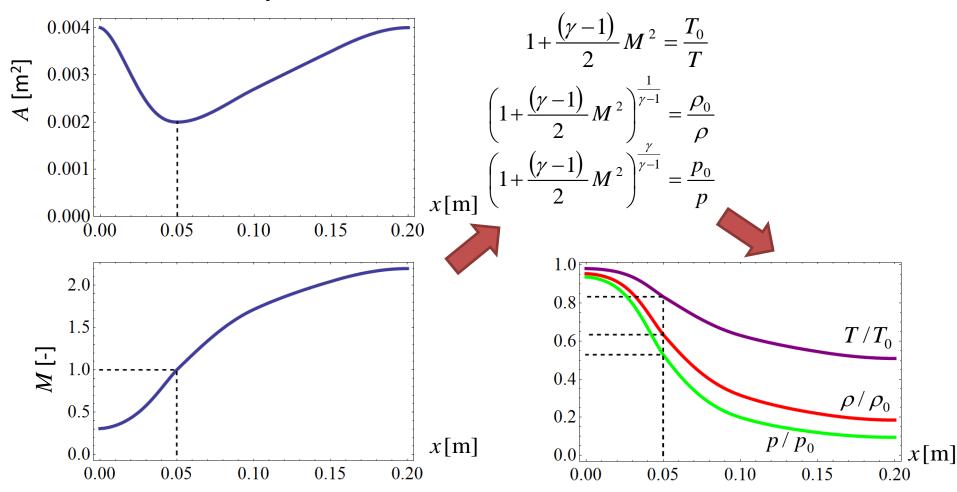
Nozzle fluid flow

variation of parameters in nozzle



Nozzle fluid flow

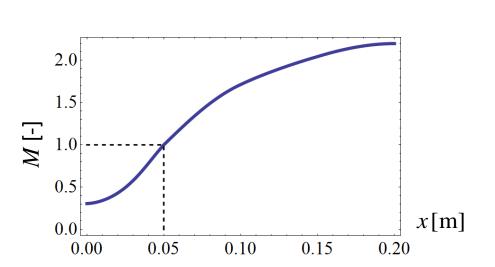
variation of parameters in nozzle

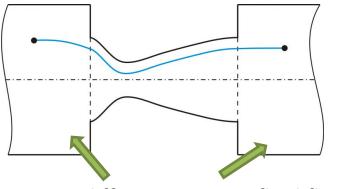


Nozzle fluid flow

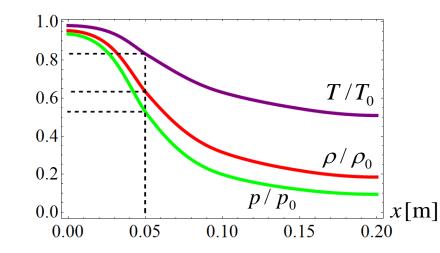
variation of parameters in nozzle

- no pressure difference no fluid flow
- if pressure ratio p_e/p_0 is different from isentropic value, the flow will be different (inside or outside the nozzle)
- exit pressure for isentropic flow with supersonic speed is p_e





pressure difference causes fluid flow



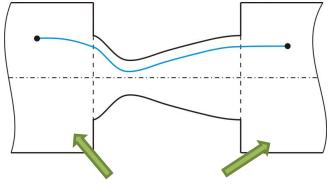
Nozzle fluid flow

variation of parameters in nozzle

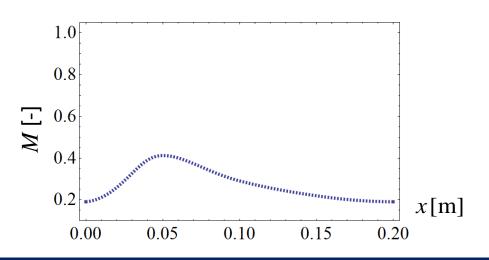
• $p_{B,1}$ is reduce below p_0

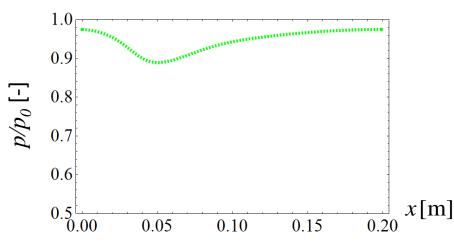


very low-speed subsonic flow, $p_{B,1} = p_{e,1}$



pressure difference causes fluid flow







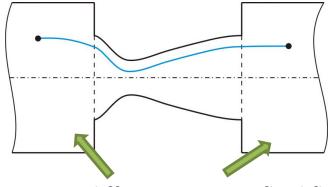
Nozzle fluid flow

variation of parameters in nozzle

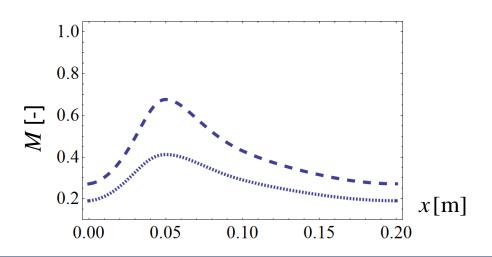
• $p_{B,2}$ is reduce below $p_{B,1}$

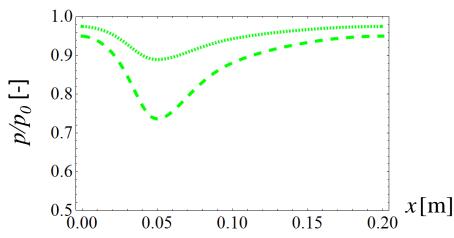


flow moves faster through nozzle, still subsonic flow, mass flow increases, $p_{B,2} = p_{e,2}$



pressure difference causes fluid flow





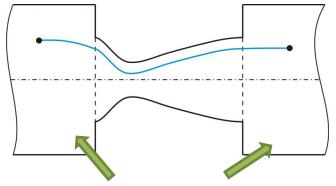
Nozzle fluid flow

variation of parameters in nozzle

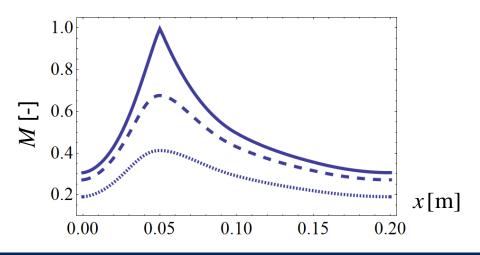
• $p_{B,3}$ is such, that it produces sonic flow in throat

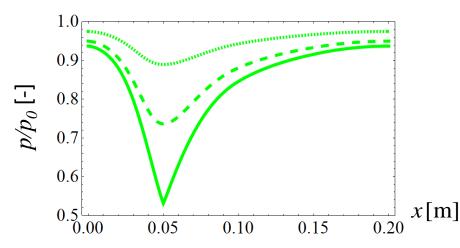


only in throat is flow sonic, in other parts of nozzle is flow subsonic, mass flow increases and reaches max. value, $p_{B.2} = p_{e.2}$



pressure difference causes fluid flow





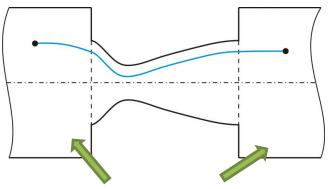
Nozzle fluid flow

variation of parameters in nozzle

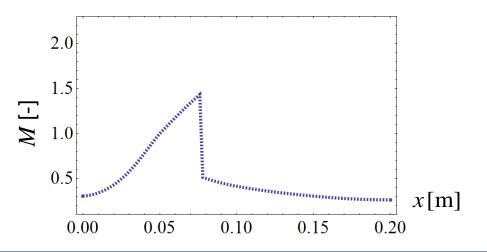
• $p_{B,4}$ is reduce below $p_{B,3}$

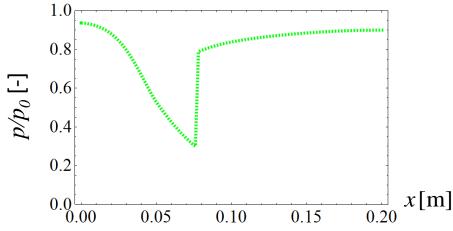


in divergent nozzle flow is at first supersonic, than shock wave is formed and flow is subsonic, mass flow is constant – chocked flow, $p_{B,4} = p_{e,4}$



pressure difference causes fluid flow





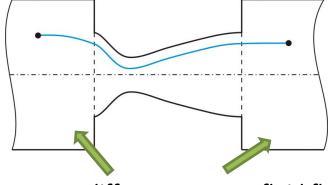
Nozzle fluid flow

variation of parameters in nozzle

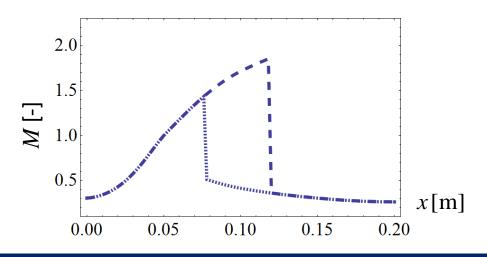
• $p_{B,5}$ is reduce below $p_{B,4}$

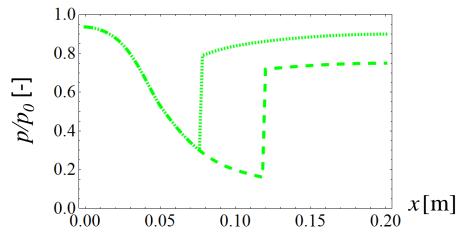


shock wave is moving toward the exit plane, $p_{B,4} = p_{e,4}$



pressure difference causes fluid flow





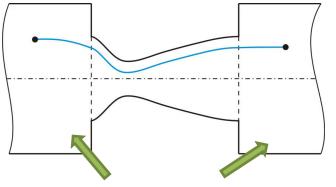
Nozzle fluid flow

variation of parameters in nozzle

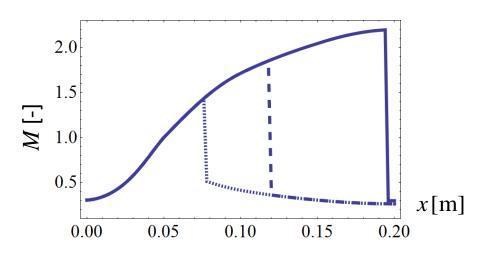
• $p_{B,6}$ is such, that shock wave is on the exit plane ___

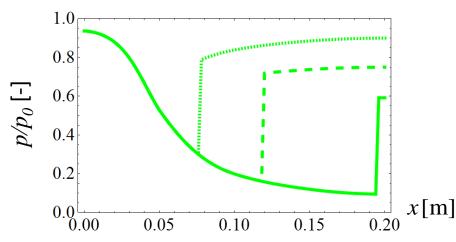


the flow is supersonic in whole nozzle except the exit plane, $p_{B,6} = p_{e,6}$



pressure difference causes fluid flow





Nozzle fluid flow

variation of parameters in nozzle

• $p_{B,6}$ is such, that shock wave is on the exit plane ___

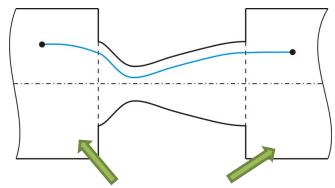


the flow is supersonic in whole nozzle except the exit plane, $p_{B.6} = p_{e.6}$

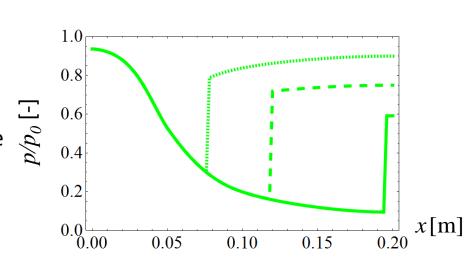


other reduction of back pressure p_B :

- exit pressure is constant p_e
- if $p_B > p_e$ shock waves moves outside nozzle
- if $p_B = p_e$ no shock waves are produced
- if $p_B < p_e$ expansion waves are formed outside the nozzle



pressure difference causes fluid flow





Nozzle fluid flow

variation of parameters in nozzle

over-expanded flow $p_B > p_e$ (low altitudes)

• $p_{B,6}$ is such, that shock wave is on the exit plane ___



the flow is supersonic in whole nozzle except the exit plane, $p_{B,6} = p_{e,6}$



other reduction of back pressure p_B :

- exit pressure is constant p_e
- if $p_B > p_e$ shock waves moves outside nozzle
- if $p_B = p_e$ no shock waves are produced
- if $p_B < p_e$ expansion waves are formed outside the nozzle





under-expanded flow $p_B < p_e$ (high altitudes)

source: NASA



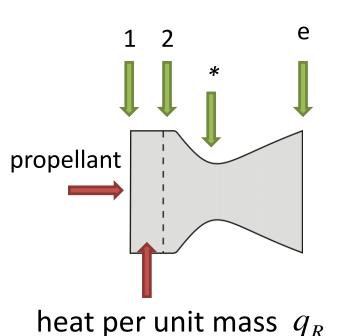


- Performance Characteristics
- Liquid Propellant Performance
- Feed System

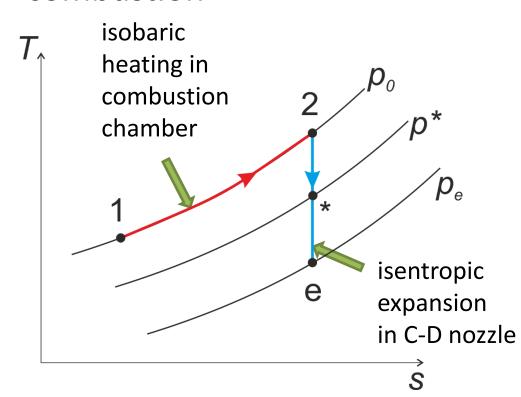


Performance Characteristics

Rocket thrust chamber – combustion



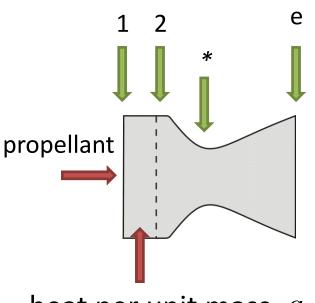
First thermodynamic law – isobaric heating:



$$q_R = \Delta h = c_p \Delta T$$

Performance Characteristics

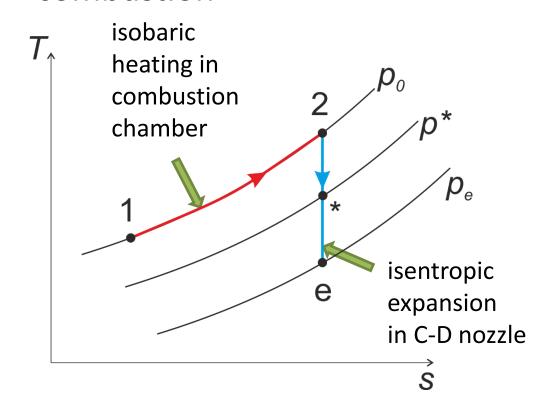
Rocket thrust chamber – combustion



heat per unit mass q_R



First thermodynamic law – isobaric heating:

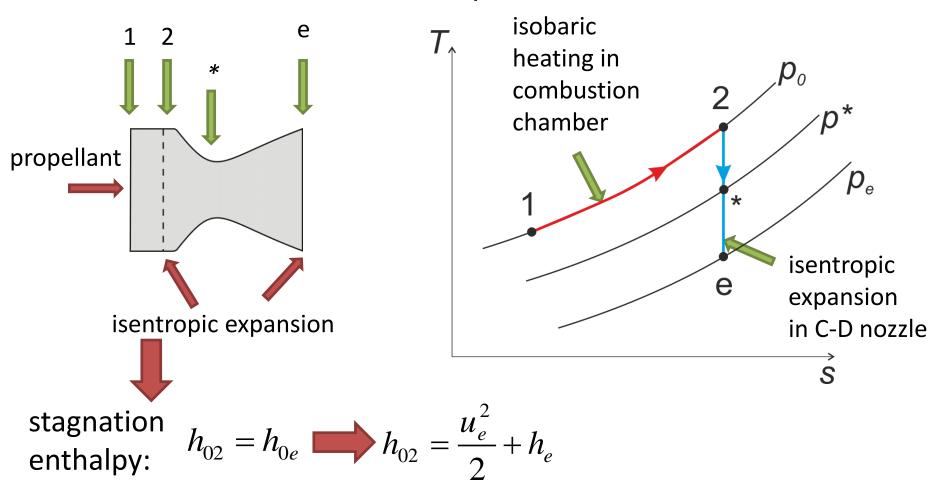


$$q_R = \Delta h = c_p \Delta T$$

$$T_{02} = T_{01} + \frac{q_R}{c_n}$$

Performance Characteristics

Rocket thrust chamber – expansion

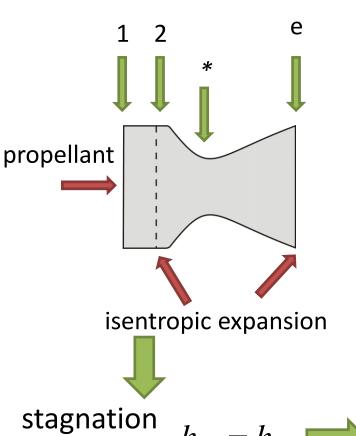


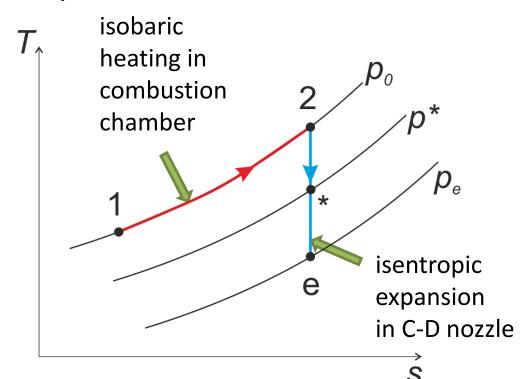




Performance Characteristics

Rocket thrust chamber – expansion





enthalpy:

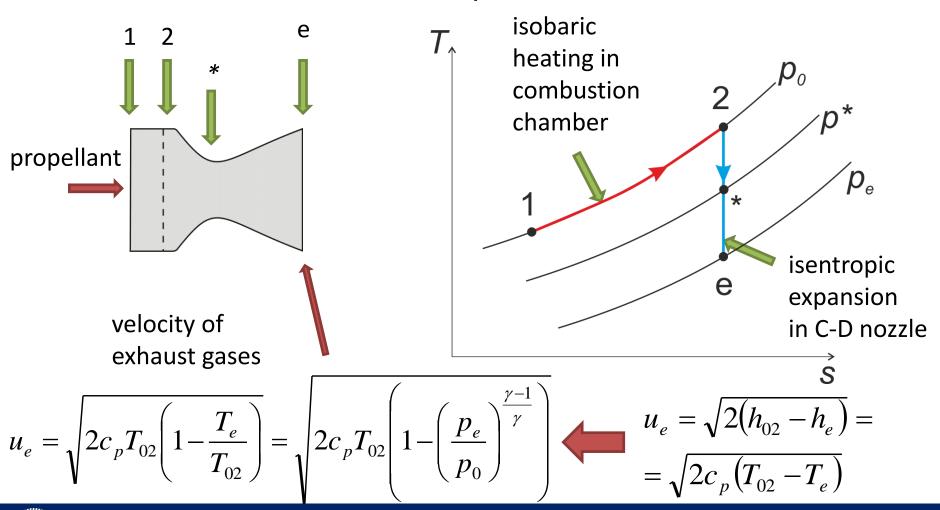
$$h_{02} = h_{0e} \implies h_{02} = \frac{u_e^2}{2} + h_e \implies \frac{u_e^2}{2}$$

$$u_e = \sqrt{2(h_{02} - h_e)} = \sqrt{2c_p(T_{02} - T_e)}$$



Performance Characteristics

Rocket thrust chamber – expansion

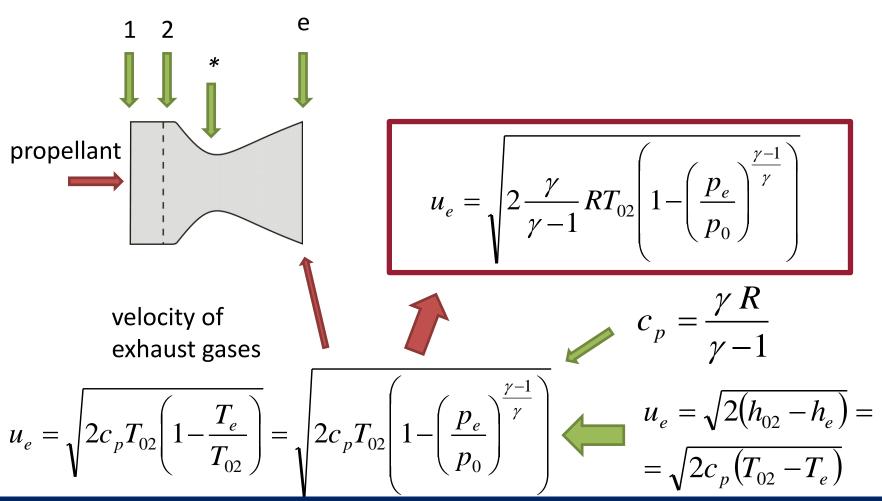






Performance Characteristics

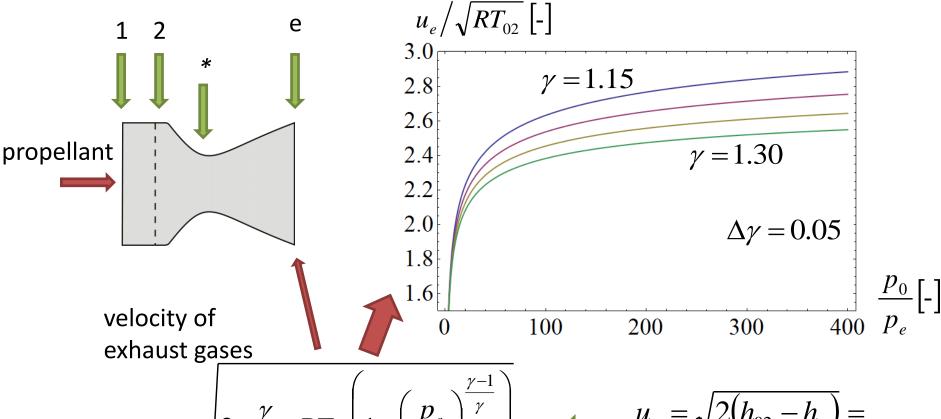
Rocket thrust chamber – exit velocity



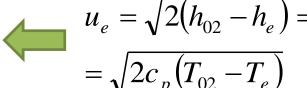


Performance Characteristics

Rocket thrust chamber – exit velocity



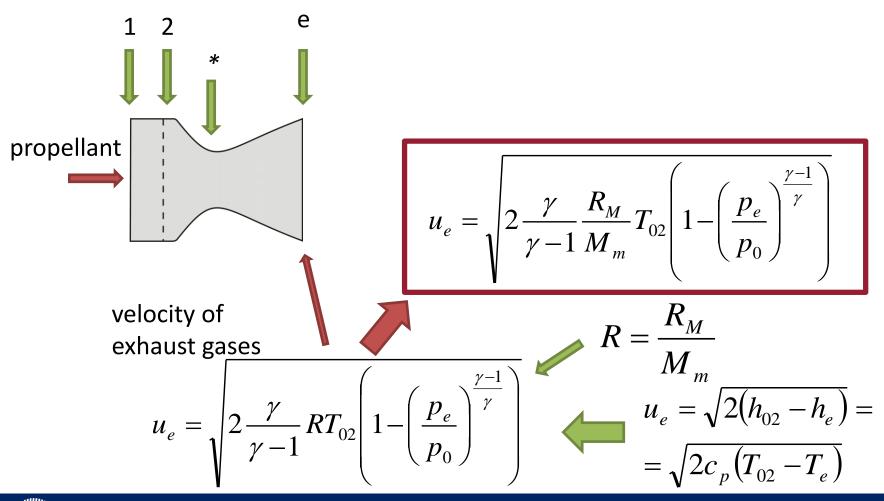
$$u_e = \sqrt{2 \frac{\gamma}{\gamma - 1} RT_{02} \left(1 - \left(\frac{p_e}{p_0} \right)^{\frac{\gamma - 1}{\gamma}} \right)}$$





Performance Characteristics

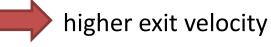
Rocket thrust chamber – exit velocity



Performance Characteristics



higher stag. temperature



low molecular weight



higher exit velocity

higher pressure ratio p_0/p_e

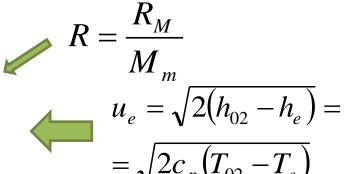


higher exit velocity

$$u_e = \sqrt{2 \frac{\gamma}{\gamma - 1} \frac{R_M}{M_m} T_{02} \left(1 - \left(\frac{p_e}{p_0} \right)^{\frac{\gamma - 1}{\gamma}} \right)}$$

velocity of exhaust gases

$$u_{e} = \sqrt{2 \frac{\gamma}{\gamma - 1} RT_{02} \left(1 - \left(\frac{p_{e}}{p_{0}}\right)^{\frac{\gamma - 1}{\gamma}}\right)}$$

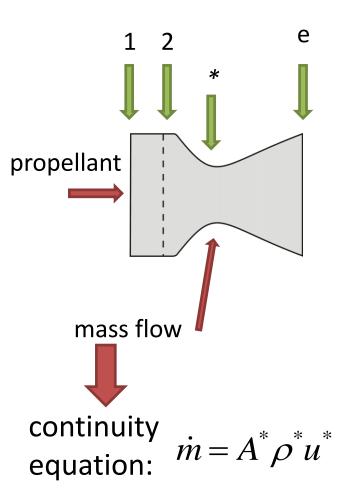




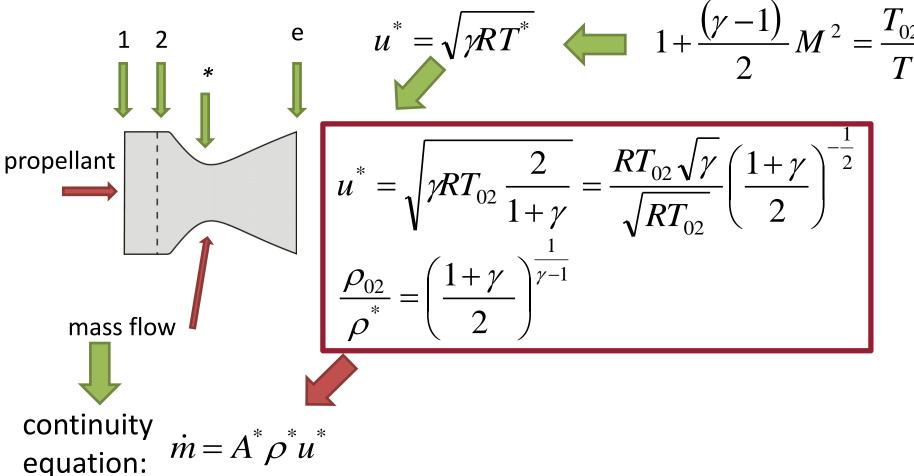
propellant



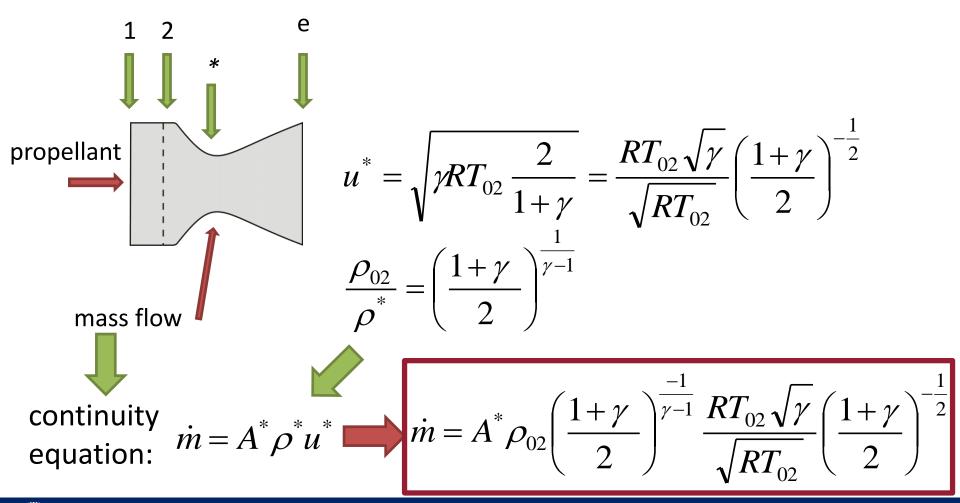
Performance Characteristics



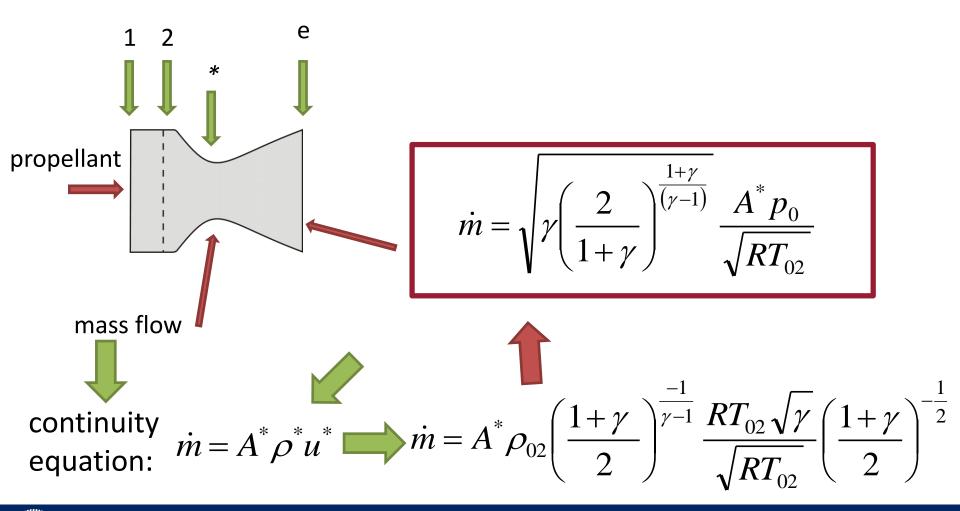
Performance Characteristics



Performance Characteristics

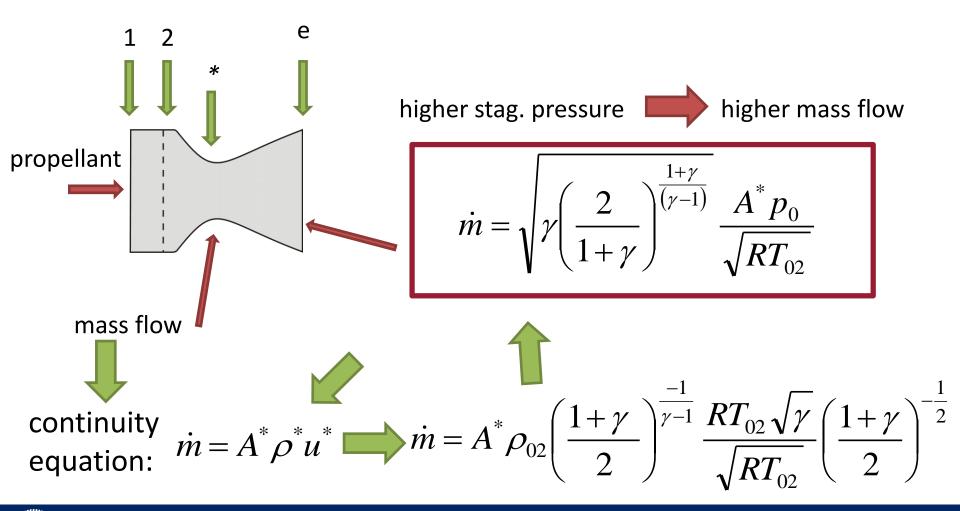


Performance Characteristics



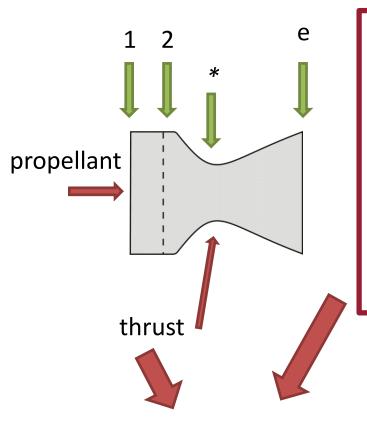


Performance Characteristics



Performance Characteristics

Rocket thrust chamber – thrust



$$u_{e} = \sqrt{2 \frac{\gamma}{\gamma - 1} RT_{02} \left(1 - \left(\frac{p_{e}}{p_{0}}\right)^{\frac{\gamma - 1}{\gamma}}\right)}$$

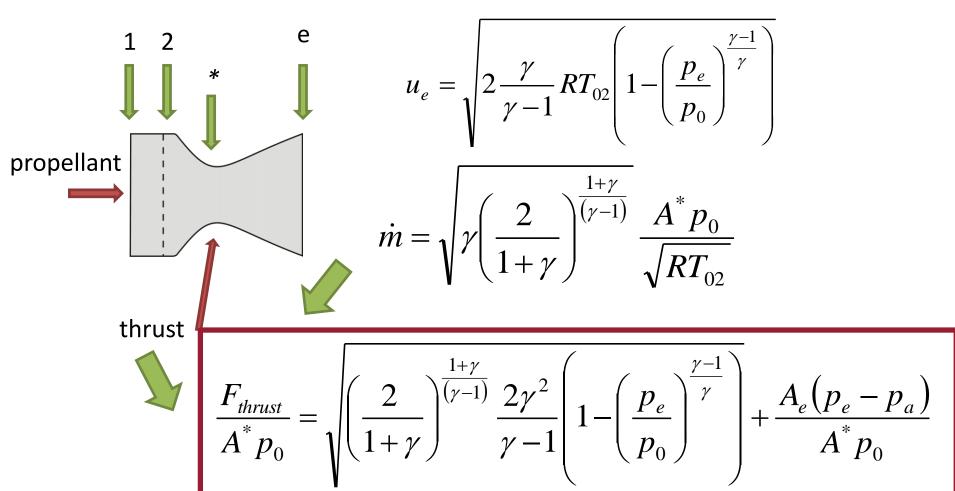
$$\frac{1 + \gamma}{\sqrt{\frac{1 + \gamma}{\gamma}}}$$

$$\dot{m} = \sqrt{\gamma \left(\frac{2}{1+\gamma}\right)^{\frac{1+\gamma}{(\gamma-1)}}} \frac{A^* p_0}{\sqrt{RT_{02}}}$$

$$F_{thrust} = \dot{m}u_e + A_e (p_e - p_a)$$

Performance Characteristics

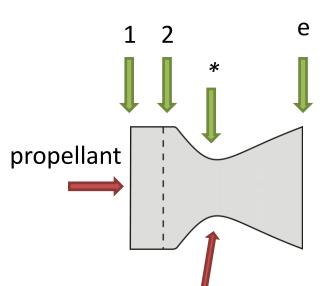
Rocket thrust chamber – thrust





Performance Characteristics

Rocket thrust chamber – thrust



$$F_{thrust} = \dot{m} rac{p_0 A^*}{\dot{m}} rac{F_{thrust}}{p_0 A^*} =$$

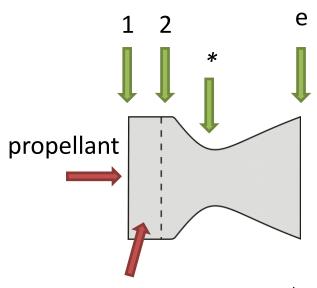
$$= \dot{m} c^* C_F$$
characteristic thrust coefficient

thrust depends only on stagnation pressure in combustion chamber

$$\frac{F_{thrust}}{A^* p_0} = \sqrt{\left(\frac{2}{1+\gamma}\right)^{\frac{1+\gamma}{(\gamma-1)}}} \frac{2\gamma^2}{\gamma - 1} \left(1 - \left(\frac{p_e}{p_0}\right)^{\frac{\gamma-1}{\gamma}}\right) + \frac{A_e(p_e - p_a)}{A^* p_0}$$

Performance Characteristics

Rocket thrust chamber – characteristic velocity



characteristic velocity – c^* specify combustion chamber

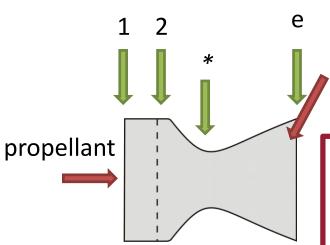
$$c^* = \frac{p_0 A^*}{\dot{m}}$$

characteristic velocity is function of combustion chamber design and propellant characteristics

$$c^* = \sqrt{\frac{1}{\gamma} \left(\frac{\gamma + 1}{2}\right)^{\frac{\gamma + 1}{\gamma - 1}} \frac{R_M T_{02}}{M_m}}$$

Performance Characteristics

Rocket thrust chamber – thrust coefficient



thrust coefficient
$$C_F$$

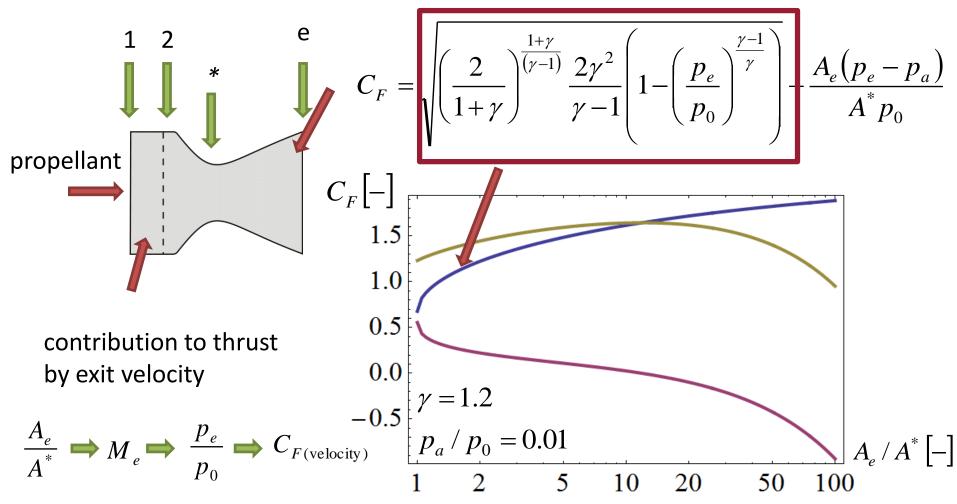
$$C_F = \frac{F_{thrust}}{p_0 A^*}$$

$$C_{F} = \sqrt{\left(\frac{2}{1+\gamma}\right)^{\frac{1+\gamma}{(\gamma-1)}} \frac{2\gamma^{2}}{\gamma-1} \left(1 - \left(\frac{p_{e}}{p_{0}}\right)^{\frac{\gamma-1}{\gamma}}\right)} + \frac{A_{e}(p_{e} - p_{a})}{A^{*}p_{0}}$$

thrust coefficient depends on gas property (γ) and nozzle parameters (nozzle area ratio and pressure ratio), it is independent on combustion chamber temperature

Performance Characteristics

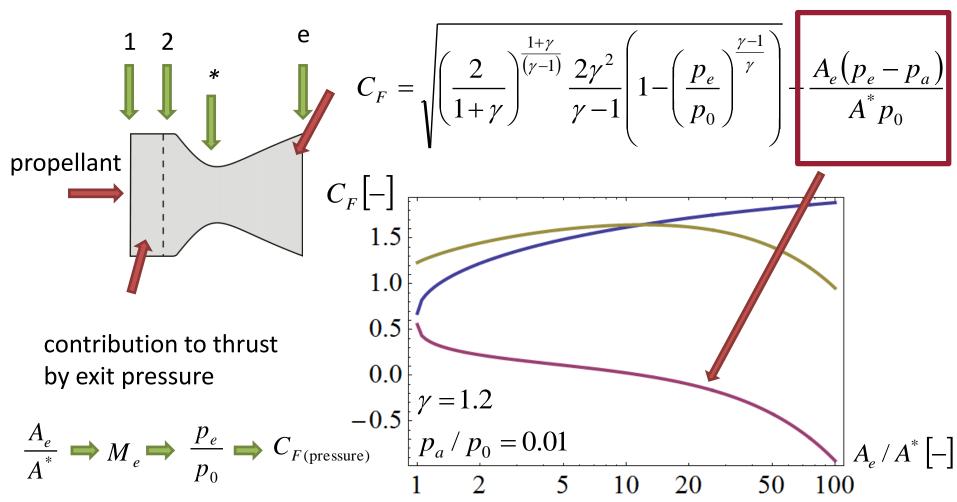
Rocket thrust chamber – thrust coefficient





Performance Characteristics

Rocket thrust chamber – thrust coefficient

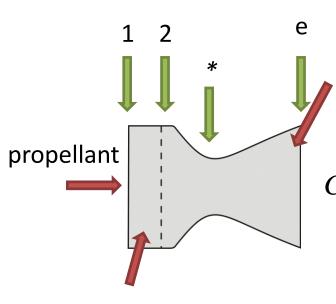




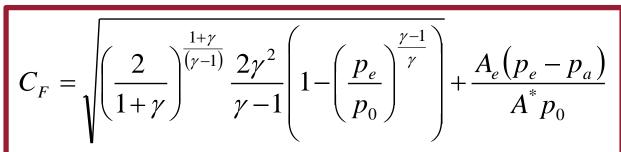


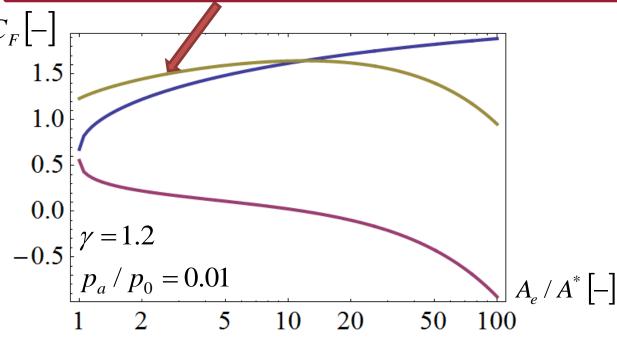
Performance Characteristics

Rocket thrust chamber – thrust coefficient



thrust coefficient for defined conditions $\left.p_a\right/p_0$ depends on area ratio



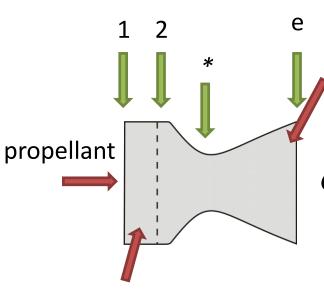




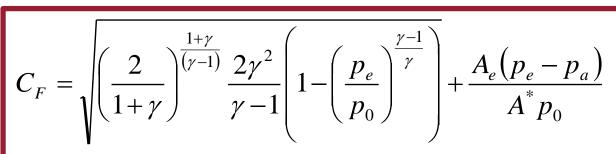


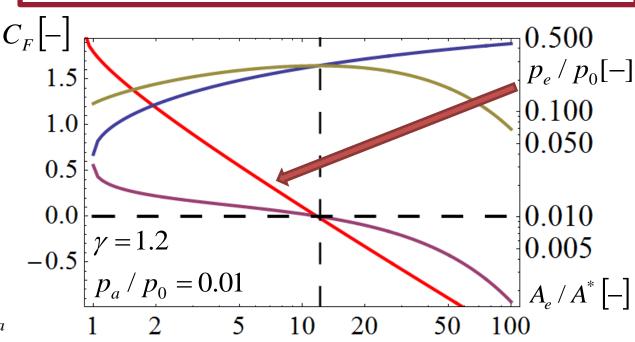
Performance Characteristics

Rocket thrust chamber – thrust coefficient



optimal thrust coefficient is defined by area ratio, where exit pressure equals ambient pressure $p_e = p_a$



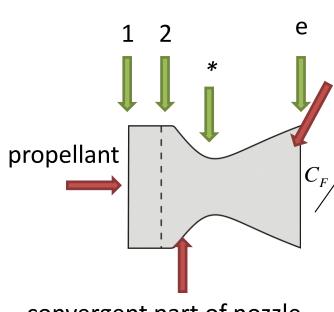






Performance Characteristics

Rocket thrust chamber – thrust coefficient



$$C_{F} = \sqrt{\left(\frac{2}{1+\gamma}\right)^{\frac{1+\gamma}{(\gamma-1)}} \frac{2\gamma^{2}}{\gamma-1} \left(1 - \left(\frac{p_{e}}{p_{0}}\right)^{\frac{\gamma-1}{\gamma}}\right)} + \frac{A_{e}(p_{e} - p_{a})}{A^{*}p_{0}}$$

convergent part of nozzle

$$M = 1$$

$$\left(1 + \frac{(\gamma - 1)}{2}M^{2}\right)^{\frac{\gamma}{\gamma - 1}} = \frac{p_{0}}{p_{e}}$$

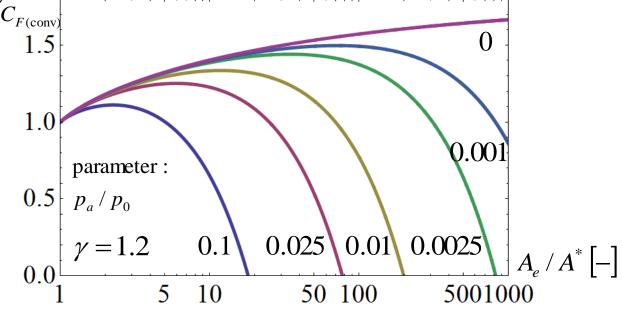
$$C_{F(\text{conv})} \leftarrow C_{F}$$

$$0.5$$

$$p_{a} / p_{0}$$

$$\gamma = 1.2$$

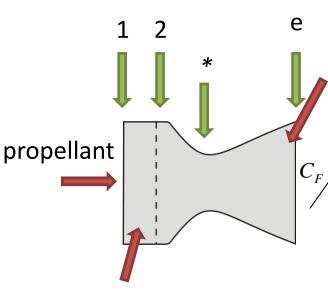
$$0.0$$



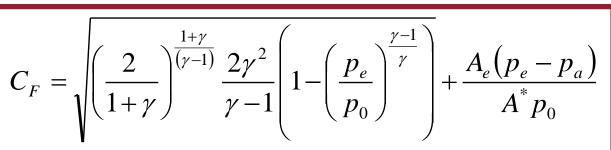


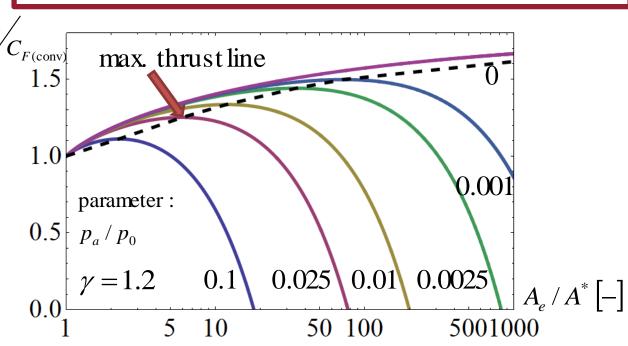
Performance Characteristics

Rocket thrust chamber – thrust coefficient



optimal thrust coefficient for individual pressure ratio form max. thrust line

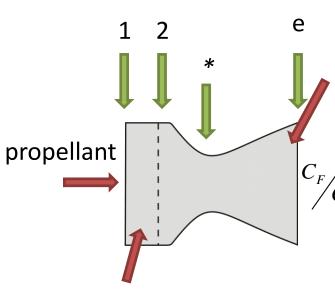




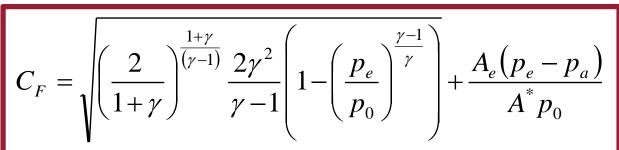


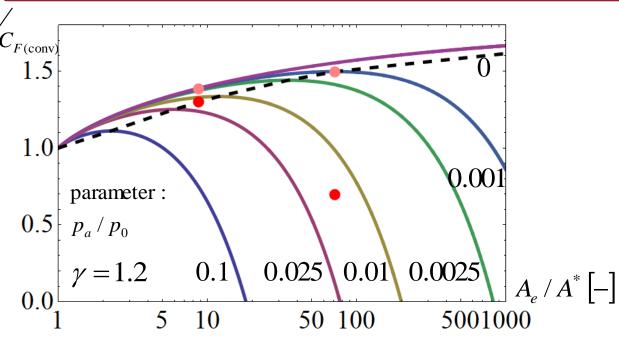
Performance Characteristics

Rocket thrust chamber – thrust coefficient



optimal thrust coefficient for sea level and pressure ratio 0.001 and the change of the coefficient

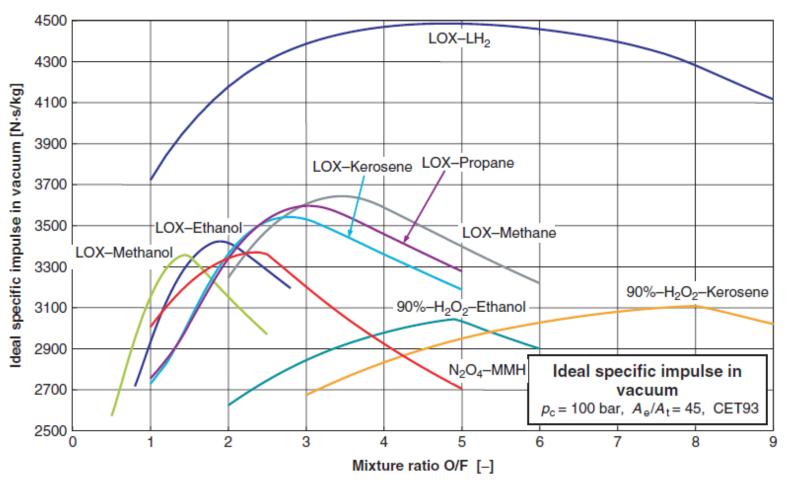






Liquid Propellant Performance

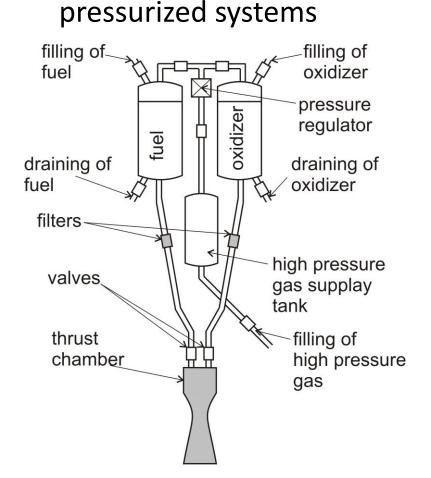
performance of individual liquid propellants



source: Ley, Wittmann, Hallmann: Handbook of space technology

Feed System

There are 2 main feed systems for liquid propellant:

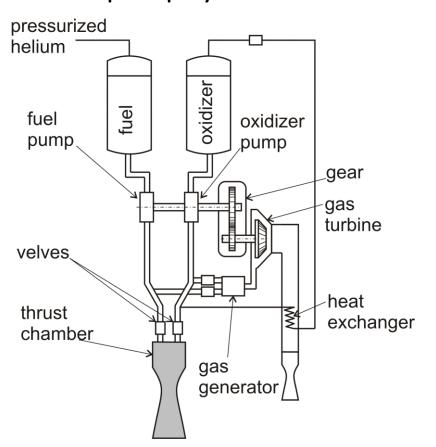


- They are usually used when:
 - total impulse is small
 - pressure in combustion chamber is small
- Disadvantages:
 - walls of tanks are thicker system is heavier
- Usage:
 - control of attitude and change of orbit



Feed System

 There are 2 main feed systems for liquid propellant: turbopump systems

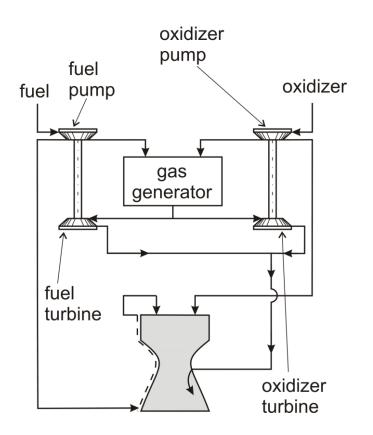


- They are usually used when:
 - total impulse is large
 - pressure in combustion chamber is large
- Positive characteristics of system:
 - pressure in tanks is lower than pressure in tanks when gas pressure feed system is used so the thickness of walls of tank is smaller
- Usage:
 - dominantly for boosters

Feed System

turbopump systems – 3 basic cycles

Gas generator cycle - open cycle



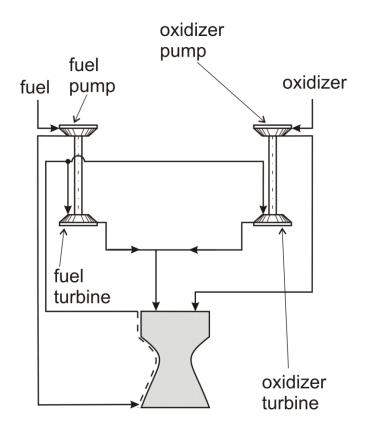
Description:

- It is the most common cycle
- It is relatively simple cycle
- The cycle efficiency is smaller than efficiency of closed cycle
- Small part of the propellant is consumed in small combustion chamber for generating gas for a turbine, which drives the pump
- Gas from turbine flows to separate nozzle or to the end part of the main nozzle, where it operates as cooler of nozzle
- Engines: F-1 (Saturn V), 2 Vulcain (Ariane 5)

Feed System

turbopump systems – 3 basic cycles

Expander cycle - closed cycle



Description:

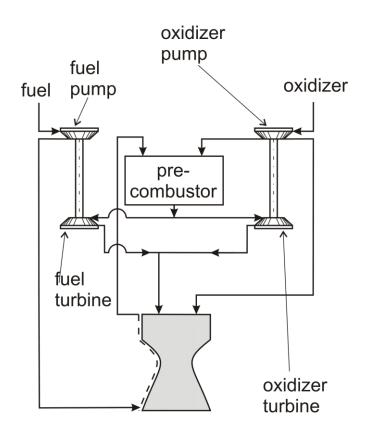
- The fuel passed through the cooling jacket of nozzle where it picked up energy and the fuel works as coolant of nozzle
- The fuel is evaporated, heated, and then fed to low pressure-ratio turbines
- at the outlet of the turbine fuel enters the combustion chamber where it is mixed with an oxidizer
- in that cycle all the fuel is burnt in combustion chamber and the efficiency of engine is increased
- Engines: RL10 (the second stage of the Delta IV),
 Vinci (ESA)



Feed System

turbopump systems – 3 basic cycles

Staged-combustion cycle - closed cycle



Description:

- The fuel passed through the cooling jacket of nozzle as in expander cycle
- Then the fuel flows into the precombustor where all the fuel is burnt with a part of the oxidizer, forming a high-energy gas to drive
- The turbines that drive the pumps all the gas at the outlet of the turbine flows into the combustion chamber where is mixed with remaining oxidizer
- pressure in combustion chamber: up to 40 MPa
- Engines: Space Shuttle Main Engine SSME, RD-170 (Energija)

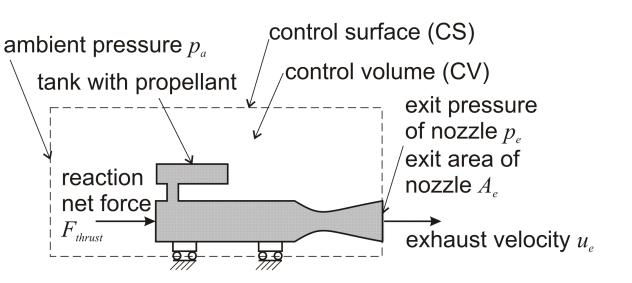


- Static Performance
- Force-Free Motion
- Motion with Gravity
- Launch Flight Mechanics



Static Performance

- Momentum equation written for CV:
- simplifications:
 - quasi-one dimensional flow
 - steady-state flow

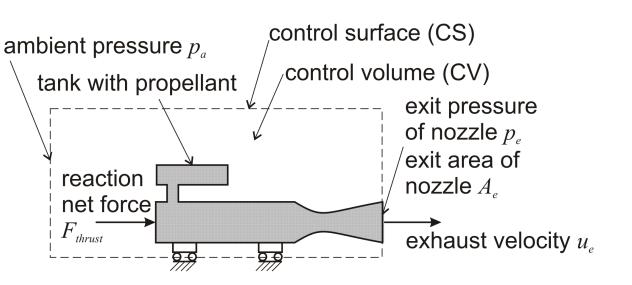






Static Performance

$$\begin{bmatrix} \text{Forces on} \\ \text{gas in CV in} \\ \text{direction } x \end{bmatrix} = \begin{bmatrix} \text{Rate} \\ \text{momentum} \\ \text{leaves CV} \end{bmatrix} - \begin{bmatrix} \text{Rate} \\ \text{momentum} \\ \text{enters CV} \end{bmatrix}$$

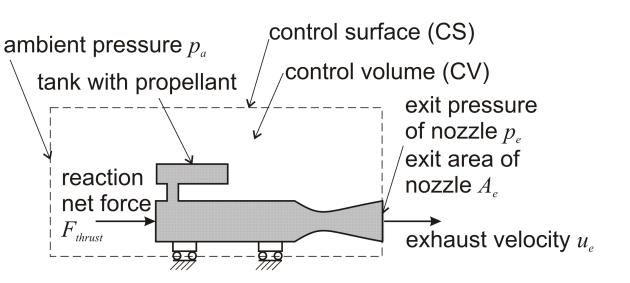






Static Performance

Forces on gas in CV in direction
$$x$$
 = Rate momentum leaves CV Rate momentum enters CV $F_{thrust} + A_e(p_a - p_e) = \dot{m}u_e$

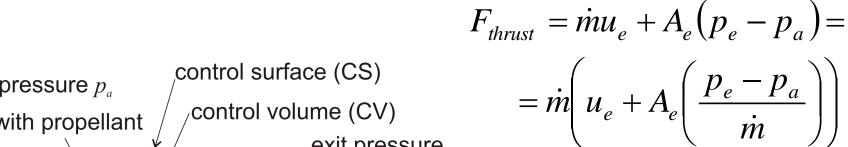


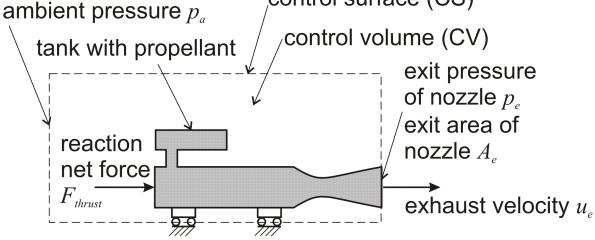




Static Performance

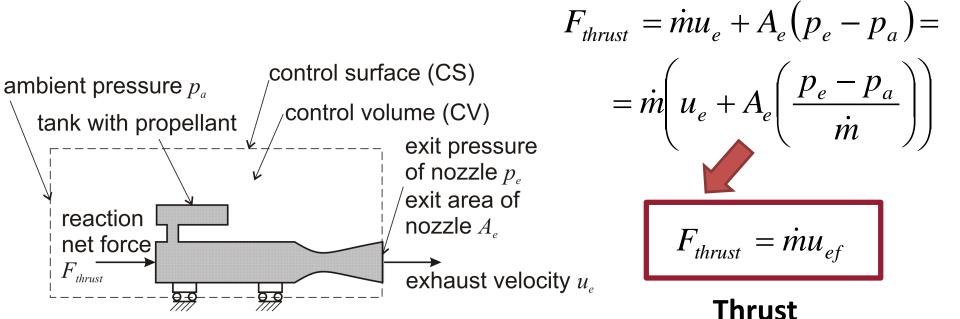
$$\begin{bmatrix} \text{Forces on} \\ \text{gas in CV in} \\ \text{direction } x \end{bmatrix} = \begin{bmatrix} \text{Rate} \\ \text{momentum} \\ \text{leaves CV} \end{bmatrix} - \begin{bmatrix} \text{Rate} \\ \text{momentum} \\ \text{enters CV} \end{bmatrix} \longrightarrow F_{thrust} + A_e (p_a - p_e) = \dot{m}u_e$$





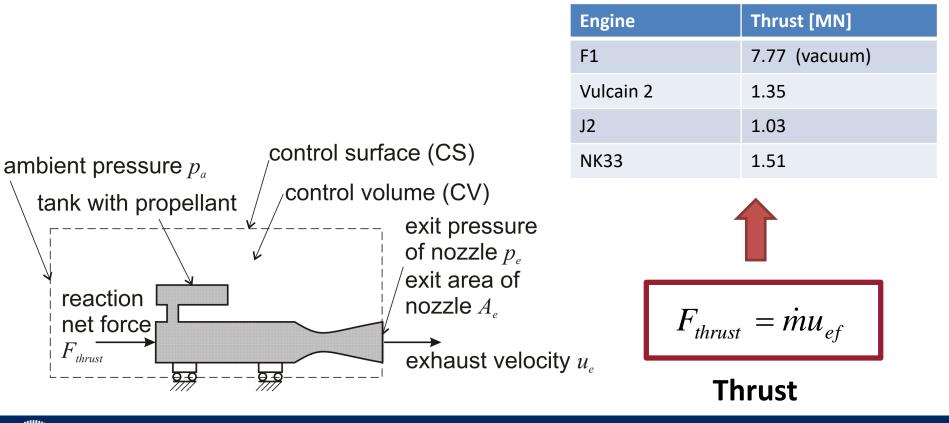
Static Performance

$$\begin{bmatrix} \text{Forces on} \\ \text{gas in CV in} \\ \text{direction } x \end{bmatrix} = \begin{bmatrix} \text{Rate} \\ \text{momentum} \\ \text{leaves CV} \end{bmatrix} - \begin{bmatrix} \text{Rate} \\ \text{momentum} \\ \text{enters CV} \end{bmatrix} \longrightarrow F_{thrust} + A_e (p_a - p_e) = \dot{m}u_e$$



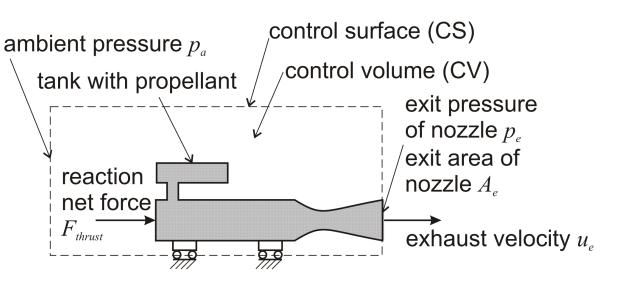


Static Performance



Static Performance

Momentum equation – written for CV:



Total impulse

$$I_{t} = F_{thrust} t$$

$$F_{thrust} = \dot{m}u_{ef}$$

Thrust





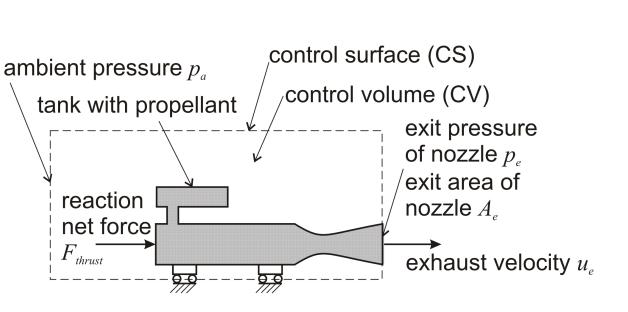
Static Performance

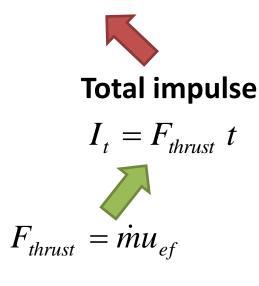
Momentum equation – written for CV:

Specific impulse

$$I_s = I_t / mg = F_{thrust} / \dot{m}g =$$

$$= \dot{m}u_{ef} / \dot{m}g = u_{ef} / g$$





Thrust

Static Performance

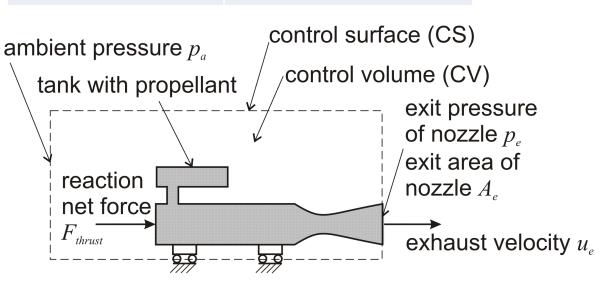
Momentum equation – written for CV:

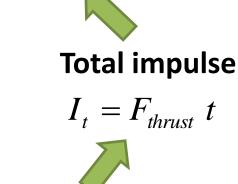
	•	
Propellant	Specific impulse I _s [s]	
cold gas	50	
Monopropellant hydrazine	230	4
LOX/LH2	455	
Ion propulsion	>3000	

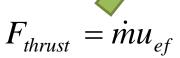


$$I_{s} = I_{t} / mg = F_{thrust} / \dot{m}g =$$

$$= \dot{m}u_{ef} / \dot{m}g = u_{ef} / g$$







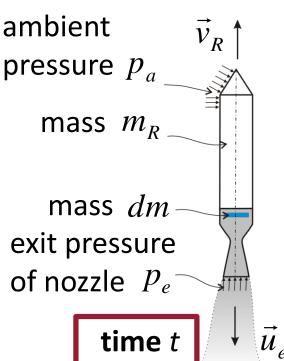
Thrust



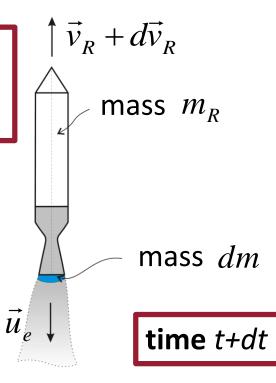


Force-Free Motion

Force-free motion absence of external forces



only ambient pressure is considered



Momentum:

$$(m_R + dm)\vec{v}_R$$



$$m_R \left(\vec{v}_R + d\vec{v}_R \right) + dm \left(\vec{v}_R + \vec{u}_e \right)$$

Change of momentum in dt

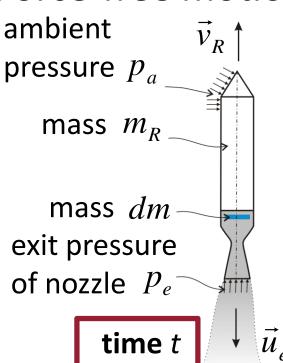
$$m_R d\vec{v}_R + dm\vec{u}_e$$



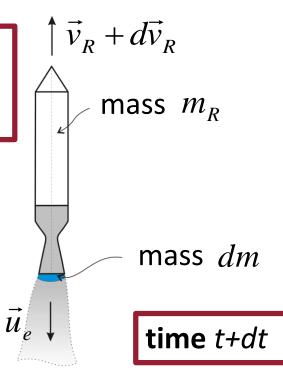


Force-Free Motion

Force-free motion absence of external forces



only ambient pressure is considered



Pressure force:
$$(p_e - p_a)A_e\vec{i}_R$$



$$(p_e - p_a)A_e \vec{i}_R$$

Total impulse in dt:

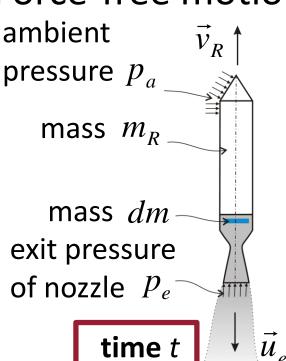
$$(p_e - p_a)A_e \vec{i}_R dt$$



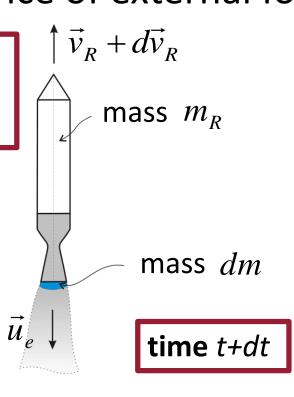


Force-Free Motion

Force-free motion absence of external forces



only ambient pressure is considered



Momentum equation: $m_R d\vec{v}_R + dm\vec{u}_o = (p_o - p_a)A_o\vec{i}_R dt$

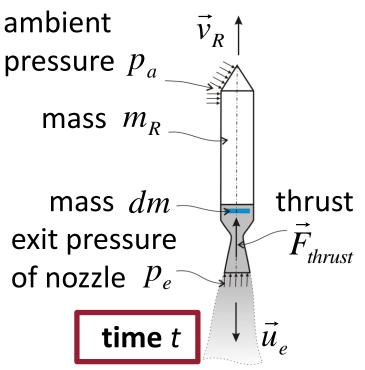
Momentum equation in \vec{i}_R : $m_R dv_R = dmu_e + (p_e - p_a)A_e dt$





Force-Free Motion

Force-free motion absence of external forces



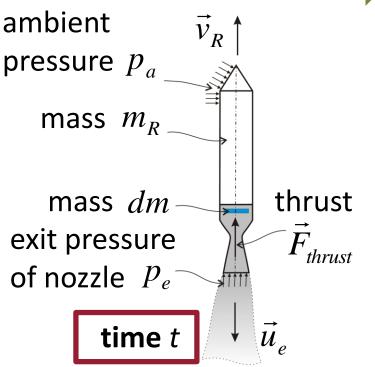
 $dm = \dot{m}dt$

Momentum equation in \vec{i}_R : $m_R dv_R = dmu_e + (p_e - p_a)A_e dt$



Force-Free Motion

Force-free motion absence of external forces



$$m_R dv_R = (\dot{m}u_e + (p_e - p_a)A_e)dt$$



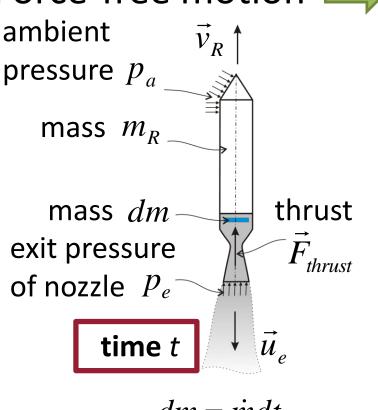
 $dm = \dot{m}dt$

Momentum equation in \vec{i}_R : $\vec{m}_R dv_R = dmu_e + (p_e - p_a)A_e dt$



Force-Free Motion

Force-free motion absence of external forces



$$dm = \dot{m}dt$$

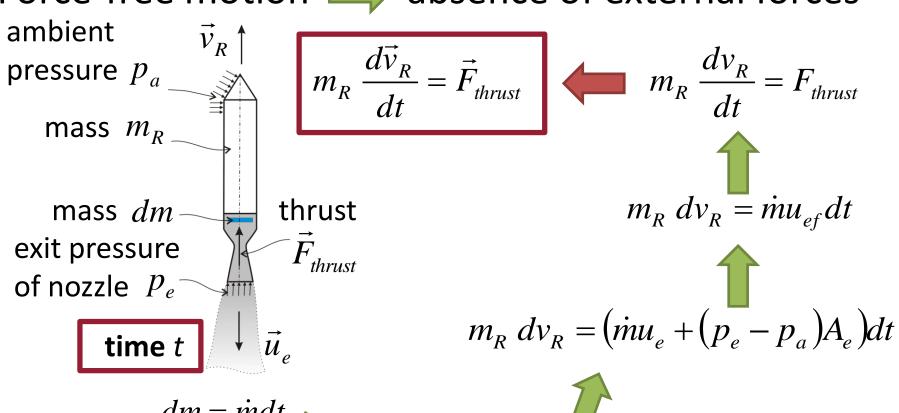
 $m_R \frac{dv_R}{dt} = F_{thrust}$ $m_R dv_R = \dot{m}u_{ef}dt$ $m_R dv_R = (\dot{m}u_e + (p_e - p_a)A_e)dt$

Momentum equation in \vec{i}_R : $\vec{m}_R dv_R = dmu_e + (p_e - p_a)A_e dt$



Force-Free Motion

absence of external forces Force-free motion



 $dm = \dot{m}dt$

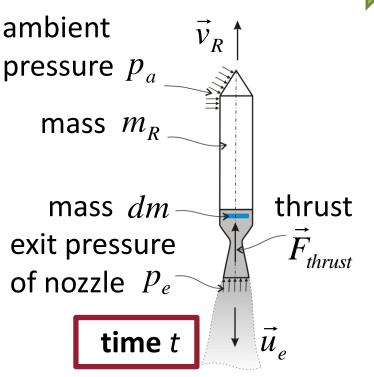
Momentum equation in i_R : $m_R dv_R = dmu_o + (p_o - p_a)A_o dt$

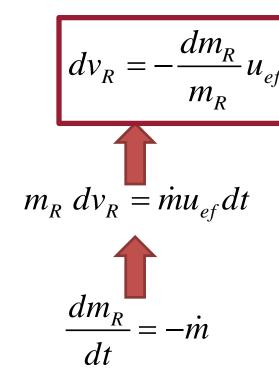




Force-Free Motion

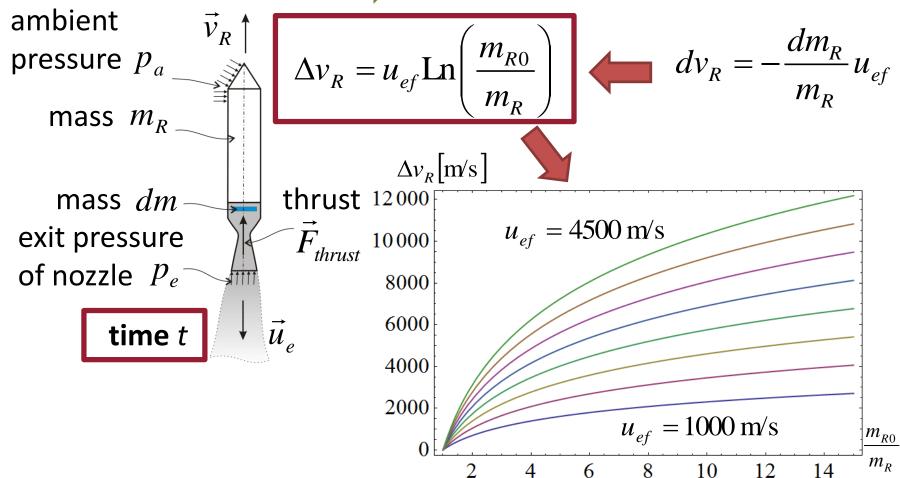
Force-free motion absence of external forces





Force-Free Motion

Force-free motion absence of external forces





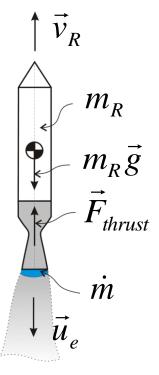


Motion with Gravity

Motion with gravity vertical motion in



gravity field



$$m_{R} \frac{dv_{R}}{dt} = -\frac{dm_{R}}{dt} u_{ef} - m_{R} g$$

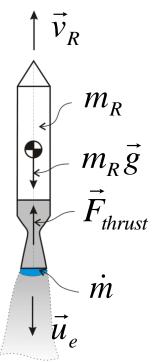
$$m_{R} \frac{d\vec{v}_{R}}{dt} = \frac{dm_{R}}{dt} \vec{u}_{ef} + m_{R} \vec{g}$$

Motion with Gravity

Motion with gravity vertical motion in



gravity field



$$dv_R = -\frac{dm_R}{m_R} u_{ef} - gdt$$

$$m_R \frac{dv_R}{dt} = -\frac{dm_R}{dt} u_{ef} - m_R g$$

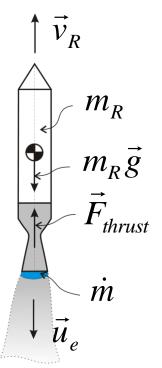
$$m_R \frac{d\vec{v}_R}{dt} = \frac{dm_R}{dt} \vec{u}_{ef} + m_R \vec{g}$$

Motion with Gravity

Motion with gravity vertical motion in



gravity field



$$dv_R = -\frac{dm_R}{m_R} u_{ef} - gdt$$

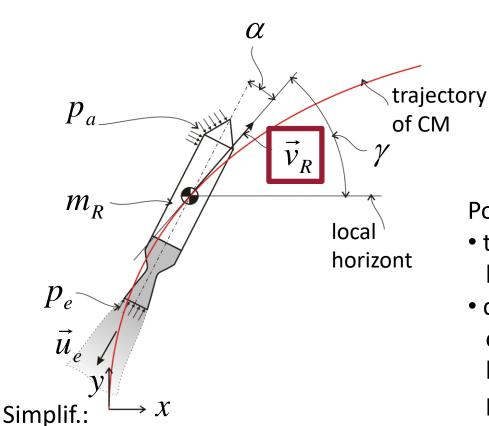


$$\Delta v_R = u_{ef} \operatorname{Ln} \left(\frac{m_{R0}}{m_R} \right) - gr$$

$$m_R \frac{dv_R}{dt} = -\frac{dm_R}{dt} u_{ef} - m_R g$$

$$m_R \frac{d\vec{v}_R}{dt} = \frac{dm_R}{dt} \vec{u}_{ef} + m_R \vec{g}$$

Forces act on rocket



- all forces act on the same plane
- Earth is inertial frame of reference

Launch Flight Mechanics

The flight of rocket has 2 main phases:

- 1. powered phase
- 2. unpowered phase

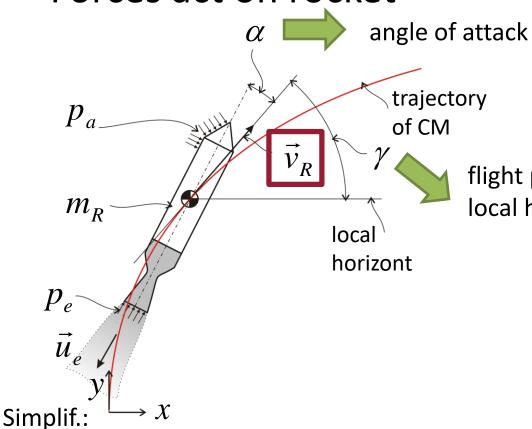


Powered phase:

- trajectory of vehicle from launch pad to burnout point
- during the phase, guidance system control the trajectory – vehicle at burnout point should have prescribed position and velocity

Launch Flight Mechanics

Forces act on rocket



flight path angle – angle between local horizont and velocity vector

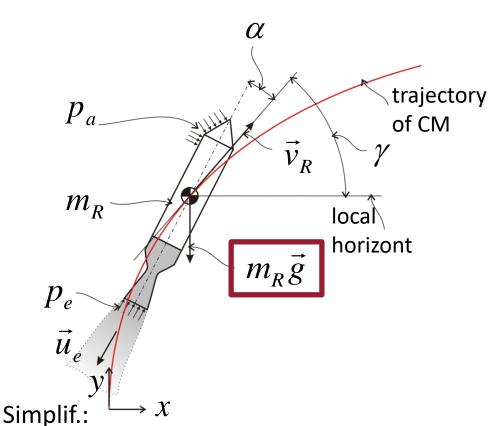
all forces act on the same plane

• Earth is inertial frame of reference





Forces act on rocket



- all forces act on the same plane
- Earth is inertial frame of reference

Launch Flight Mechanics

Three forces act on rocket at each instant:

1. gravitational force – applied at the CM $m_R \vec{g}$



it is function of vertical location of rocket mass of rocket is function of propellant mass flow

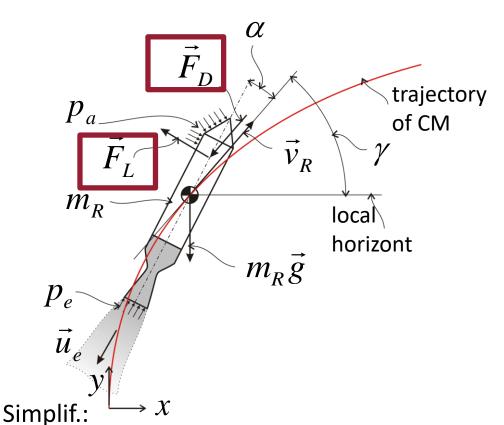


equation of propellant mass flow



Launch Flight Mechanics

Forces act on rocket



- all forces act on the same plane
- Earth is inertial frame of reference

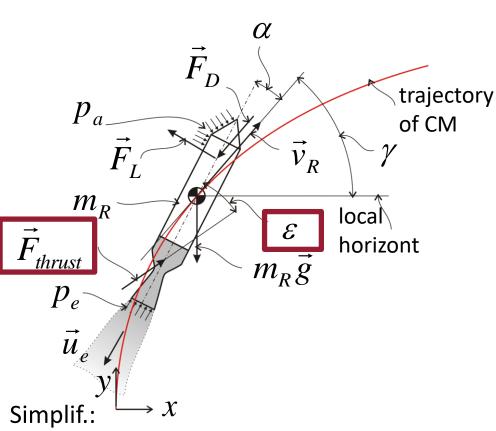
Three forces act on rocket at each instant:

- 1. gravitational force applied at the CM $m_R \vec{g}$
- 2. aerodynamic force applied at aerodynamic center and can be decomposed into:
 - drag force \vec{F}_D



they are function of vertical location and attitude of rocket

Forces act on rocket



- all forces act on the same plane
- Earth is inertial frame of reference

Launch Flight Mechanics

Three forces act on rocket at each instant:

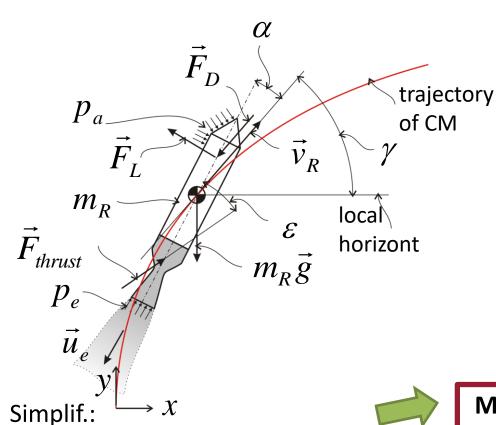
- 1. gravitational force applied at the CM $m_R \vec{g}$
- 2. aerodynamic force applied at aerodynamic center and can be decomposed into:
 - drag force \vec{F}_D
 - lift force \vec{F}_L
- 3. thrust force \vec{F}_{thrust}



magnitude and direction can be controlled

Launch Flight Mechanics

Forces act on rocket



Three forces act on rocket at each instant:

- 1. gravitational force applied at the CM $m_R \vec{g}$
- 2. aerodynamic force applied at aerodynamic center and can be decomposed into:
 - drag force \vec{F}_D
 - lift force \vec{F}_{I}
- 3. thrust force \vec{F}_{thrust}



• all forces act on the same plane

• Earth is inertial frame of reference

Motion of vehicle in 2D:

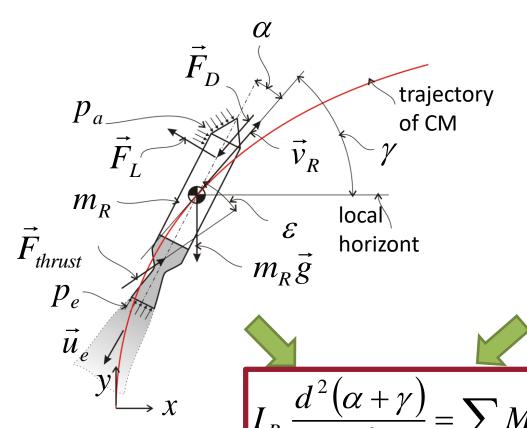
- translation motion of Center of Mass
- relative rotation motion around the CM





Launch Flight Mechanics

Forces act on rocket



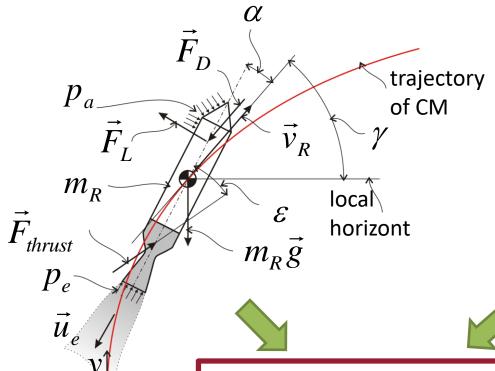
Three forces act on rocket at each instant:

- 1. gravitational force applied at the CM $m_R \vec{g}$
- 2. aerodynamic force applied at aerodynamic center and can be decomposed into:
 - drag force \vec{F}_D
 - lift force \vec{F}_{I}
- 3. thrust force \vec{F}_{thrust}

dynamic equations: relative rotation motion around the CM

Launch Flight Mechanics

Forces act on rocket



Three forces act on rocket at each instant:

- 1. gravitational force applied at the CM $m_R \vec{g}$
- 2. aerodynamic force applied at aerodynamic center and can be decomposed into:
 - drag force \vec{F}_D
- 3. thrust force \vec{F}_{thrust}

$$m_R \, \frac{d\vec{v}_R}{dt} = \vec{F}_{thrust} \, + m_R \vec{g} + \vec{F}_D + \vec{F}_L$$

dynamic equations: translation motion of the CM

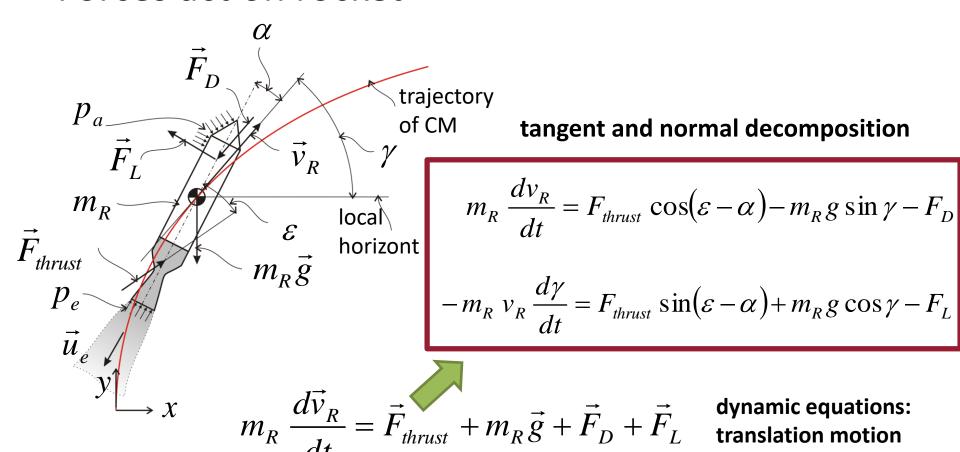




Launch Flight Mechanics

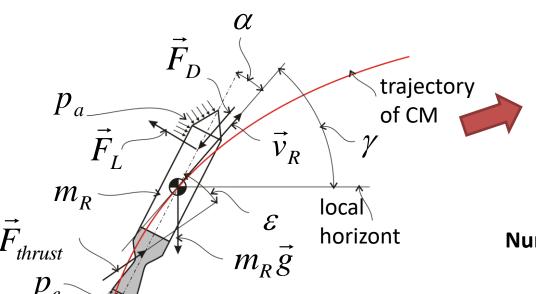
of the CM

Forces act on rocket



Launch Flight Mechanics

Forces act on rocket



Equations of rocket motion:

- dynamic equations
- kinematic equations
- equation of propellant mass flow



Numerical solution of system of ODE

$$x = \int_{0}^{t} v_R \cos \gamma \, dt$$
$$y = \int_{0}^{t} v_R \sin \gamma \, dt$$

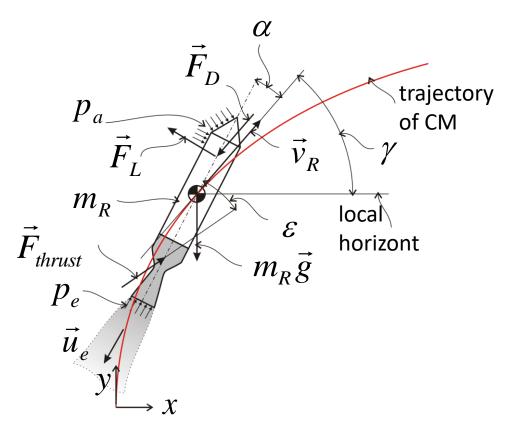
kinematic equations: vertical and horizontal distance





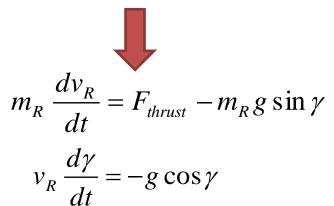
Launch Flight Mechanics

Forces act on rocket



Gravity turn:

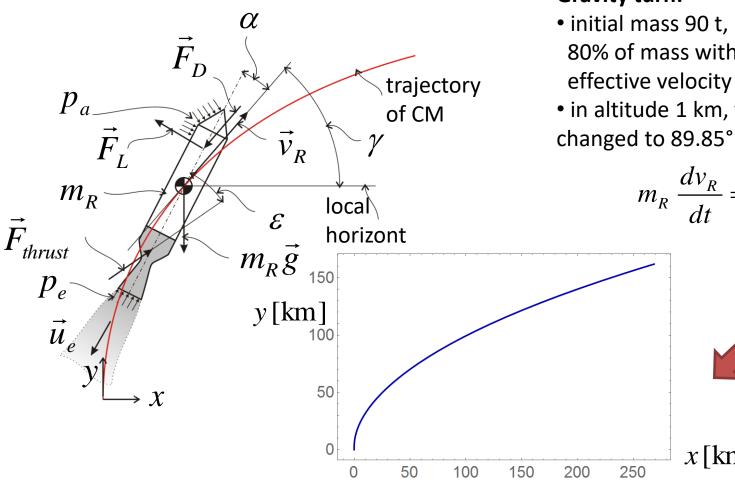
- gravity turn trajectory change of flight angle due to gravity
- only thrust and gravity is considered
- angle of attack is zero
- thrust is in axis of rocket





Launch Flight Mechanics

Forces act on rocket



Gravity turn:

- initial mass 90 t, propellant is 80% of mass with flow 250 kg/s effective velocity 4000 m/s
- in altitude 1 km, flight angle is

$$m_R \frac{dv_R}{dt} = F_{thrust} - m_R g \sin \gamma$$

$$v_R \frac{d\gamma}{dt} = -g \cos \gamma$$



$$\frac{dx}{dt} = v_R \cos \gamma$$

$$\frac{dy}{dt} = v_R \sin \gamma$$

x[km]